

# Phase diagram of the Ruby model with infinite-PEPS

**Román Orús**

*Donostia International Physics Center (DIPC)*

*January 24th 2019*

*S. S. Jahromi, RO, M. Kargarian, A. Langari, PRB 97, 115161 (2018)*

S. S. Jahromi



# Motivation: T0 with 2-body interactions

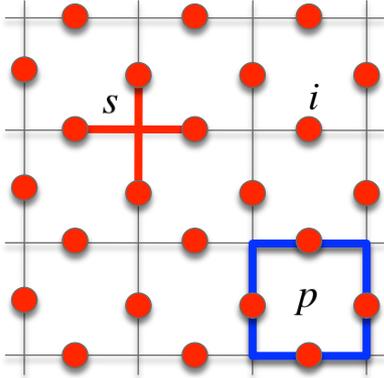


# Motivation: TQ with 2-body interactions

Toric Code

$$H = -J \sum_s A_s - J \sum_p B_p$$

$\mathbb{Z}_2$



$$A_s = \prod_{i \in s} \sigma_i^x$$

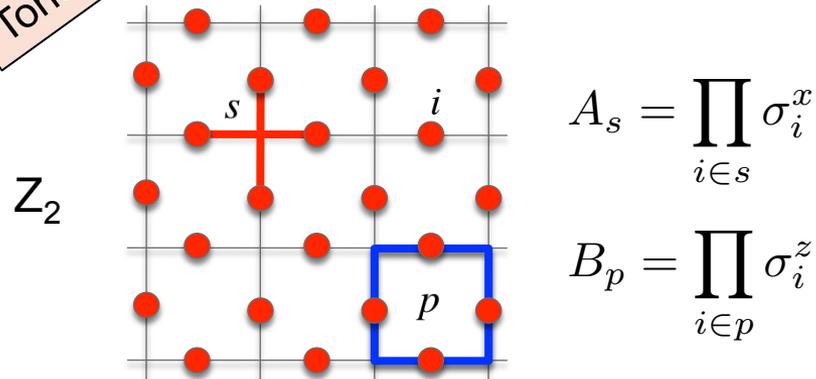
$$B_p = \prod_{i \in p} \sigma_i^z$$

A. Kitaev, *Annals of Physics* 303, 2-30 (2003)

# Motivation: TO with 2-body interactions

Toric Code

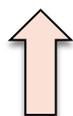
$$H = -J \sum_s A_s - J \sum_p B_p$$



$$A_s = \prod_{i \in s} \sigma_i^x$$

$$B_p = \prod_{i \in p} \sigma_i^z$$

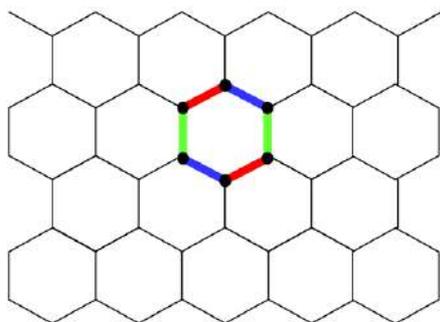
A. Kitaev, *Annals of Physics* 303, 2-30 (2003)



$$J_z \gg J_x, J_y$$

KH model

$$H = - \sum_{\alpha=x,y,z} J_\alpha \sum_{\alpha\text{-links}} \sigma_i^\alpha \sigma_j^\alpha$$



●—● YY

●—● XX

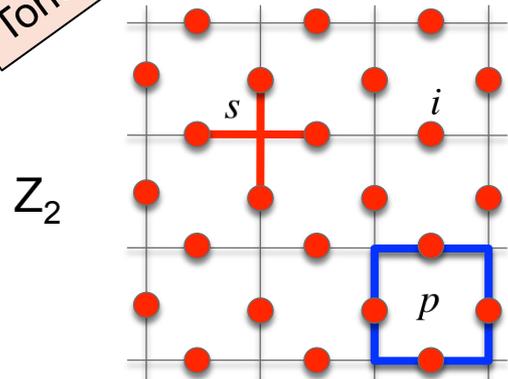
●—● ZZ

A. Kitaev, *Annals of Physics* 321, 2-111 (2006)

# Motivation: TO with 2-body interactions

Toric Code

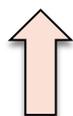
$$H = -J \sum_s A_s - J \sum_p B_p$$



$$A_s = \prod_{i \in s} \sigma_i^x$$

$$B_p = \prod_{i \in p} \sigma_i^z$$

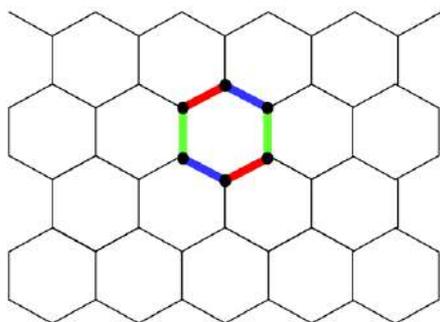
A. Kitaev, *Annals of Physics* 303, 2-30 (2003)



$$J_z \gg J_x, J_y$$

KH model

$$H = - \sum_{\alpha=x,y,z} J_\alpha \sum_{\alpha\text{-links}} \sigma_i^\alpha \sigma_j^\alpha$$



YY

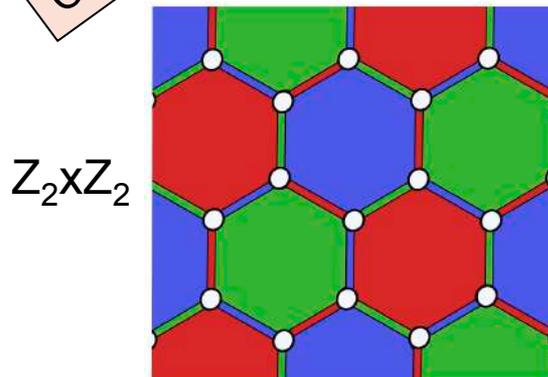
XX

ZZ

A. Kitaev, *Annals of Physics* 321, 2-111 (2006)

Color Code

$$H = -J \sum_p (X_p + Z_p)$$



$$X_p = \prod_{i \in p} \sigma_i^x$$

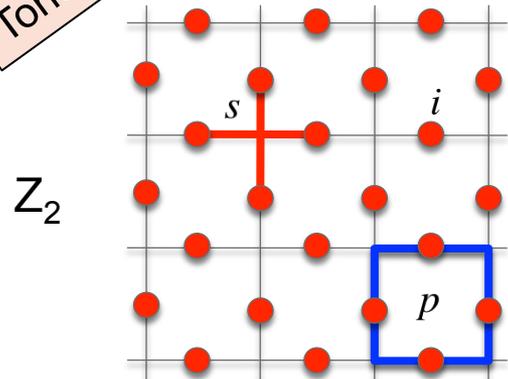
$$Z_p = \prod_{i \in p} \sigma_i^z$$

H. Bombin, M. A. Martin-Delgado, *PRL* 97 180501 (2006)

# Motivation: TO with 2-body interactions

Toric Code

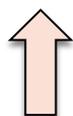
$$H = -J \sum_s A_s - J \sum_p B_p$$



$$A_s = \prod_{i \in s} \sigma_i^x$$

$$B_p = \prod_{i \in p} \sigma_i^z$$

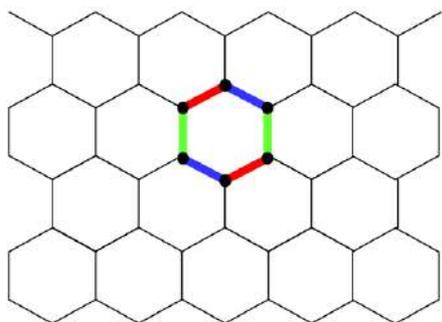
A. Kitaev, *Annals of Physics* 303, 2-30 (2003)



$$J_z \gg J_x, J_y$$

KH model

$$H = - \sum_{\alpha=x,y,z} J_\alpha \sum_{\alpha\text{-links}} \sigma_i^\alpha \sigma_j^\alpha$$



YY

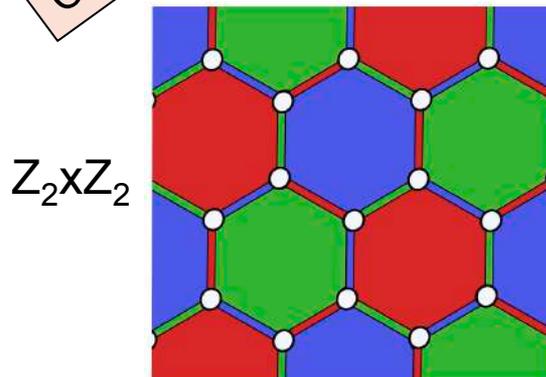
XX

ZZ

A. Kitaev, *Annals of Physics* 321, 2-111 (2006)

Color Code

$$H = -J \sum_p (X_p + Z_p)$$



$$X_p = \prod_{i \in p} \sigma_i^x$$

$$Z_p = \prod_{i \in p} \sigma_i^z$$

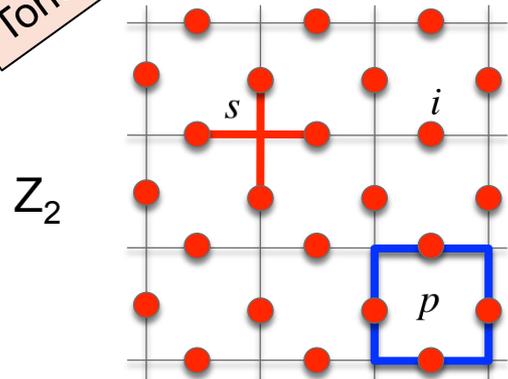
H. Bombin, M. A. Martin-Delgado, *PRL* 97 180501 (2006)



# Motivation: TO with 2-body interactions

Toric Code

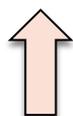
$$H = -J \sum_s A_s - J \sum_p B_p$$



$$A_s = \prod_{i \in s} \sigma_i^x$$

$$B_p = \prod_{i \in p} \sigma_i^z$$

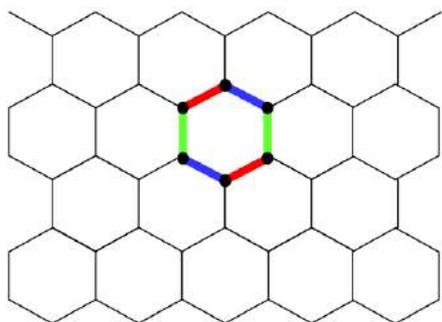
A. Kitaev, *Annals of Physics* 303, 2-30 (2003)



$$J_z \gg J_x, J_y$$

$$H = - \sum_{\alpha=x,y,z} J_\alpha \sum_{\alpha\text{-links}} \sigma_i^\alpha \sigma_j^\alpha$$

KH model



●—● YY

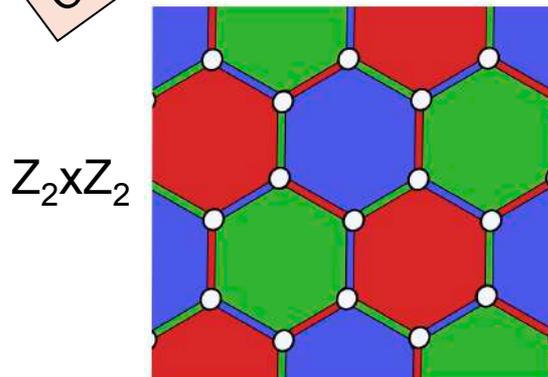
●—● XX

●—● ZZ

A. Kitaev, *Annals of Physics* 321, 2-111 (2006)

Color Code

$$H = -J \sum_p (X_p + Z_p)$$



$$X_p = \prod_{i \in p} \sigma_i^x$$

$$Z_p = \prod_{i \in p} \sigma_i^z$$

H. Bombin, M. A. Martin-Delgado, *PRL* 97 180501 (2006)

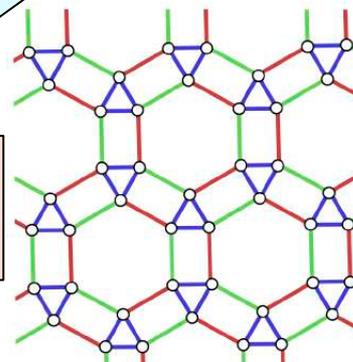


$$J_z \gg J_x, J_y$$

$$H = - \sum_z J_\alpha \sum_{\alpha\text{-links}} \sigma_i^\alpha \sigma_j^\alpha$$

Ruby model

Not exactly solvable!



●—● YY

●—● XX

●—● ZZ

S. S. Jahromi et al, *PRB* 94, 125145 (2016)

# Goal



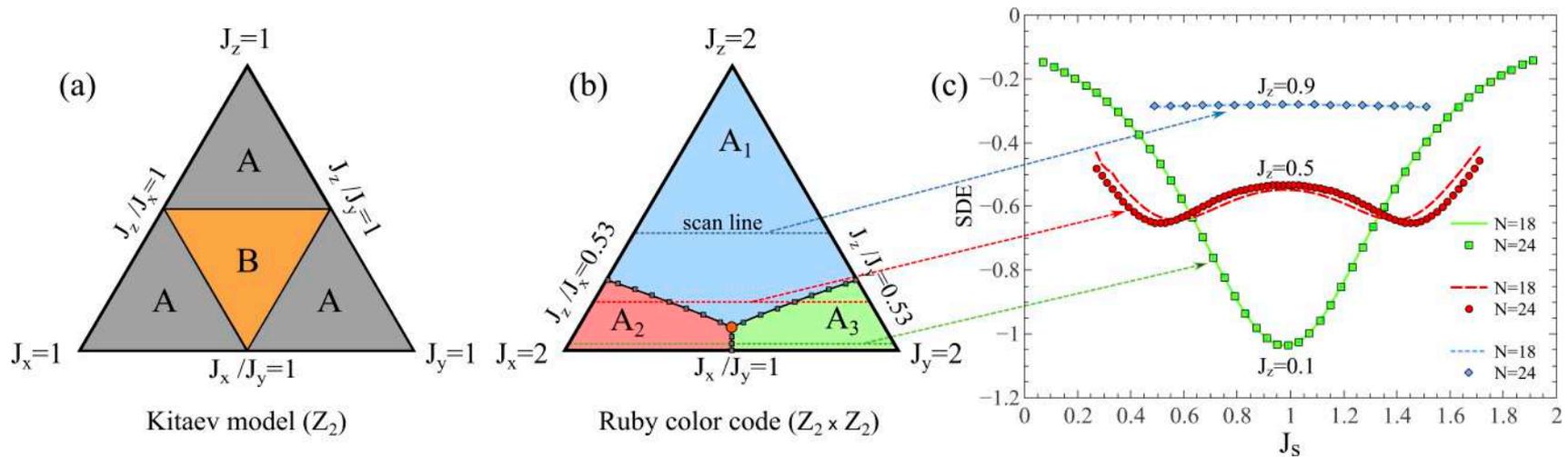
- Study the phase diagram of the ruby model with infinite-PEPS (tensor networks)

# Goal

- Study the phase diagram of the ruby model with infinite-PEPS (tensor networks)

- Previous studies by exact diagonalization (18 and 24 sites):

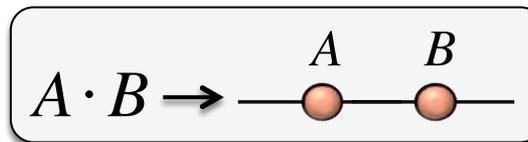
*S. S. Jahromi et al, PRB 94, 125145 (2016)*



- $A_1$  phase: gapped,  $Z_2 \times Z_2$  TO.
- $A_2$  &  $A_3$  phases: gapless, gapped by pert., conjectured Ising anyons.
- Nature of phase transitions? Accurate position?

# Tensor Networks

e.g. RO, *Annals of Physics* 349 (2014) 117–158

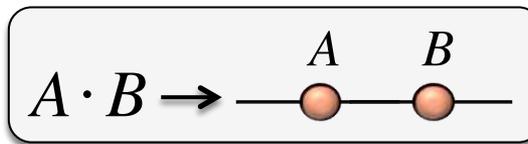


$$|\Psi\rangle = \sum_{i_s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

p-level systems

# Tensor Networks

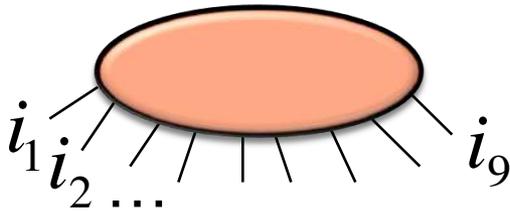
e.g. RO, *Annals of Physics* 349 (2014) 117–158



$$|\Psi\rangle = \sum_{i^s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

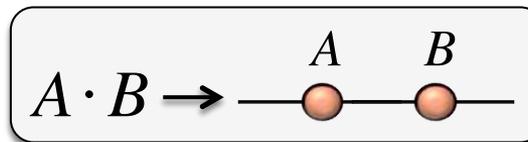
p-level systems

$$\Psi_{i_1 i_2 i_3 i_4 i_5 i_6 i_7 i_8 i_9}$$



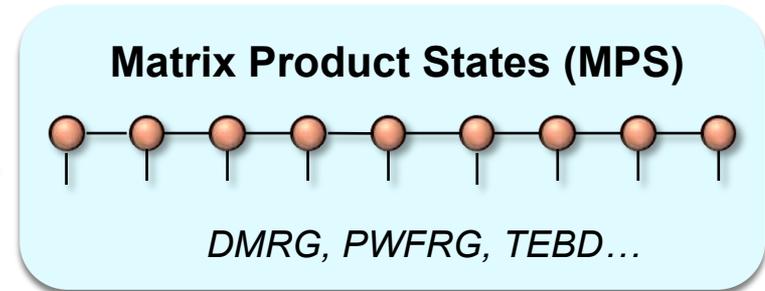
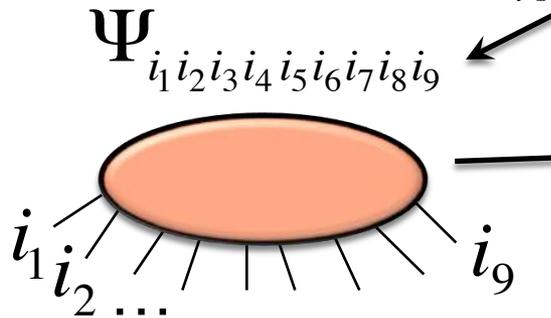
# Tensor Networks

e.g. RO, *Annals of Physics* 349 (2014) 117–158



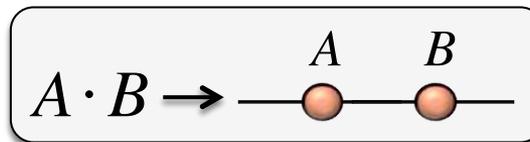
$$|\Psi\rangle = \sum_{i^s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

p-level systems



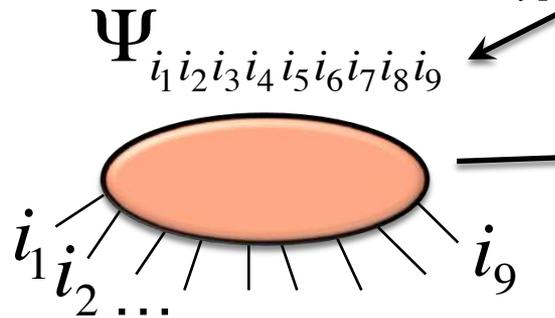
# Tensor Networks

e.g. RO, *Annals of Physics* 349 (2014) 117–158

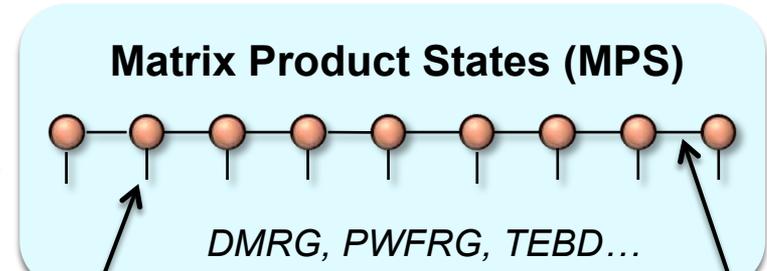


$$|\Psi\rangle = \sum_{i^s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

p-level systems



1d

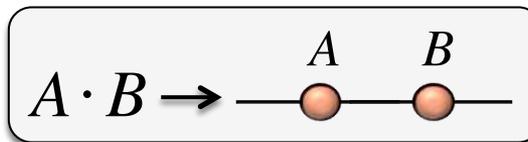


physical 1...p

bond 1..D (entanglement)

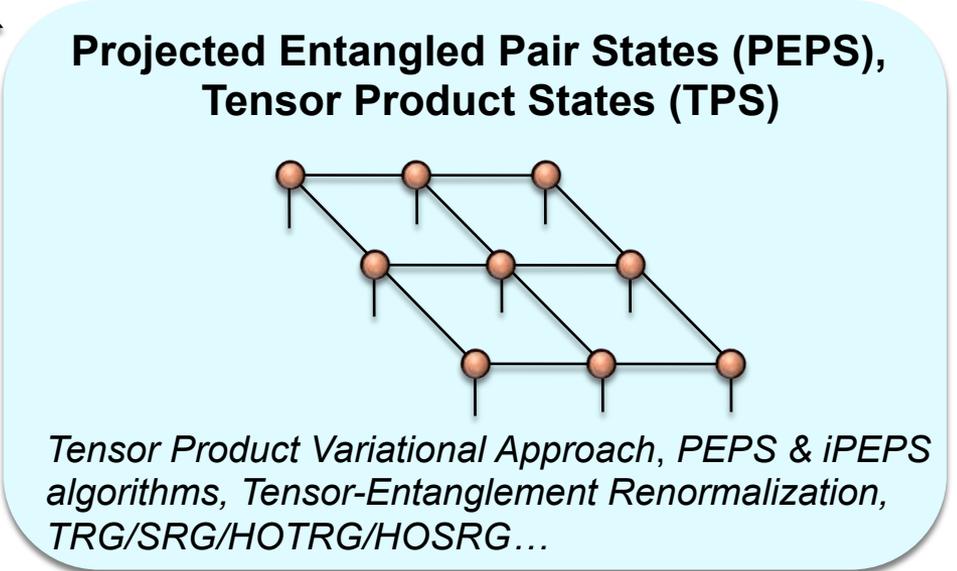
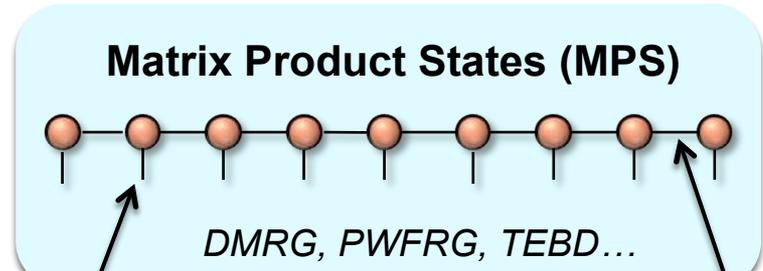
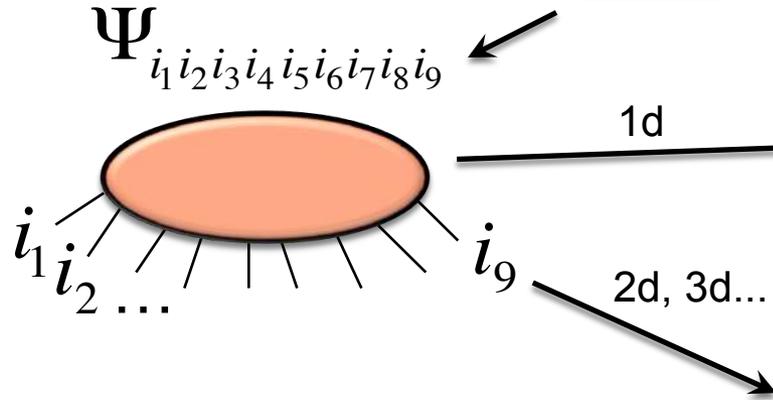
# Tensor Networks

e.g. RO, *Annals of Physics* 349 (2014) 117–158



$$|\Psi\rangle = \sum_{i^s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

p-level systems



# Tensor Networks

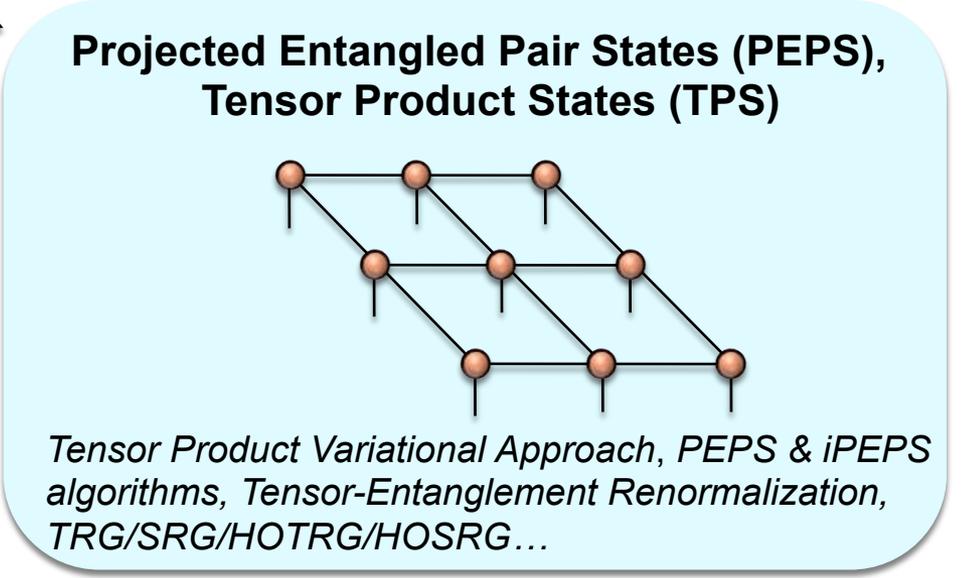
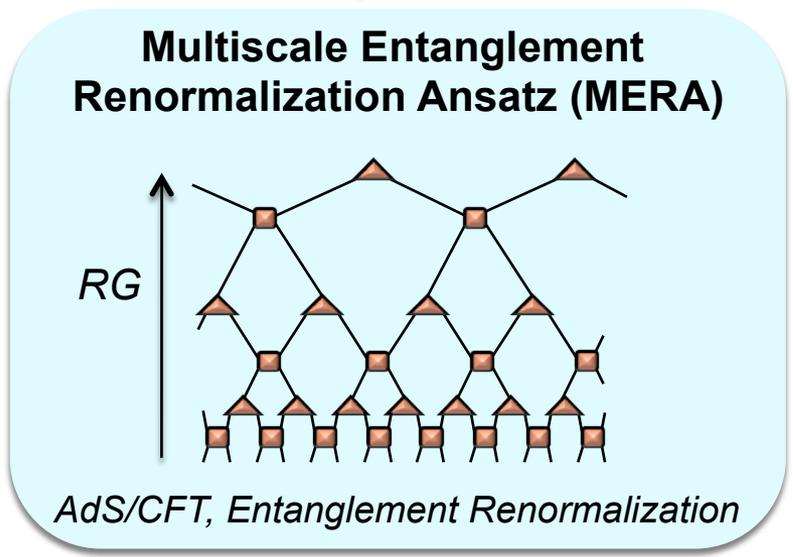
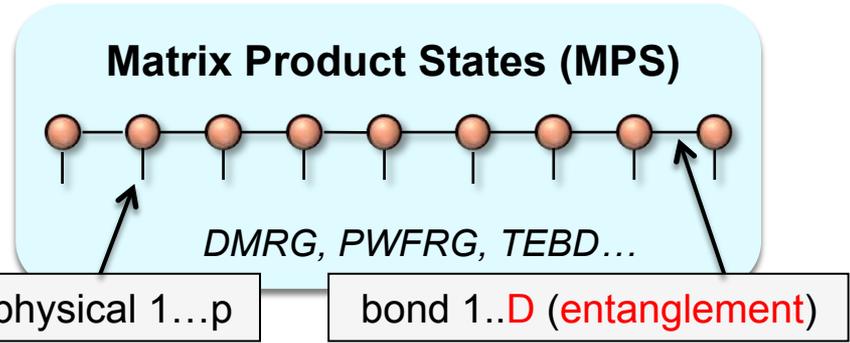
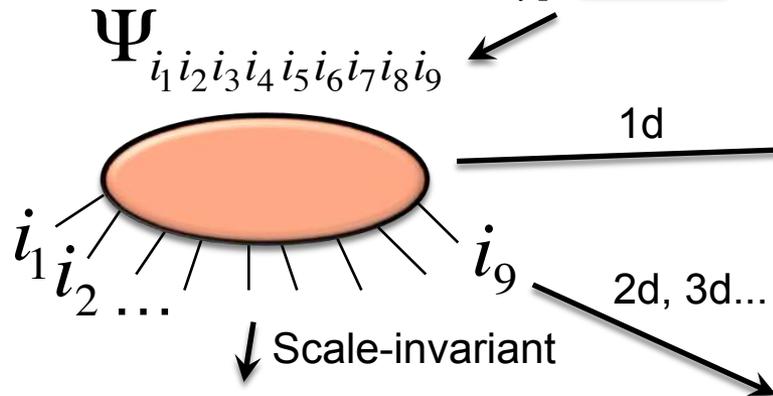
e.g. RO, *Annals of Physics* 349 (2014) 117–158



$$A \cdot B \rightarrow \text{---} \circ \text{---} \text{---} \circ \text{---}$$

$$|\Psi\rangle = \sum_{i^s} \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_N\rangle$$

p-level systems



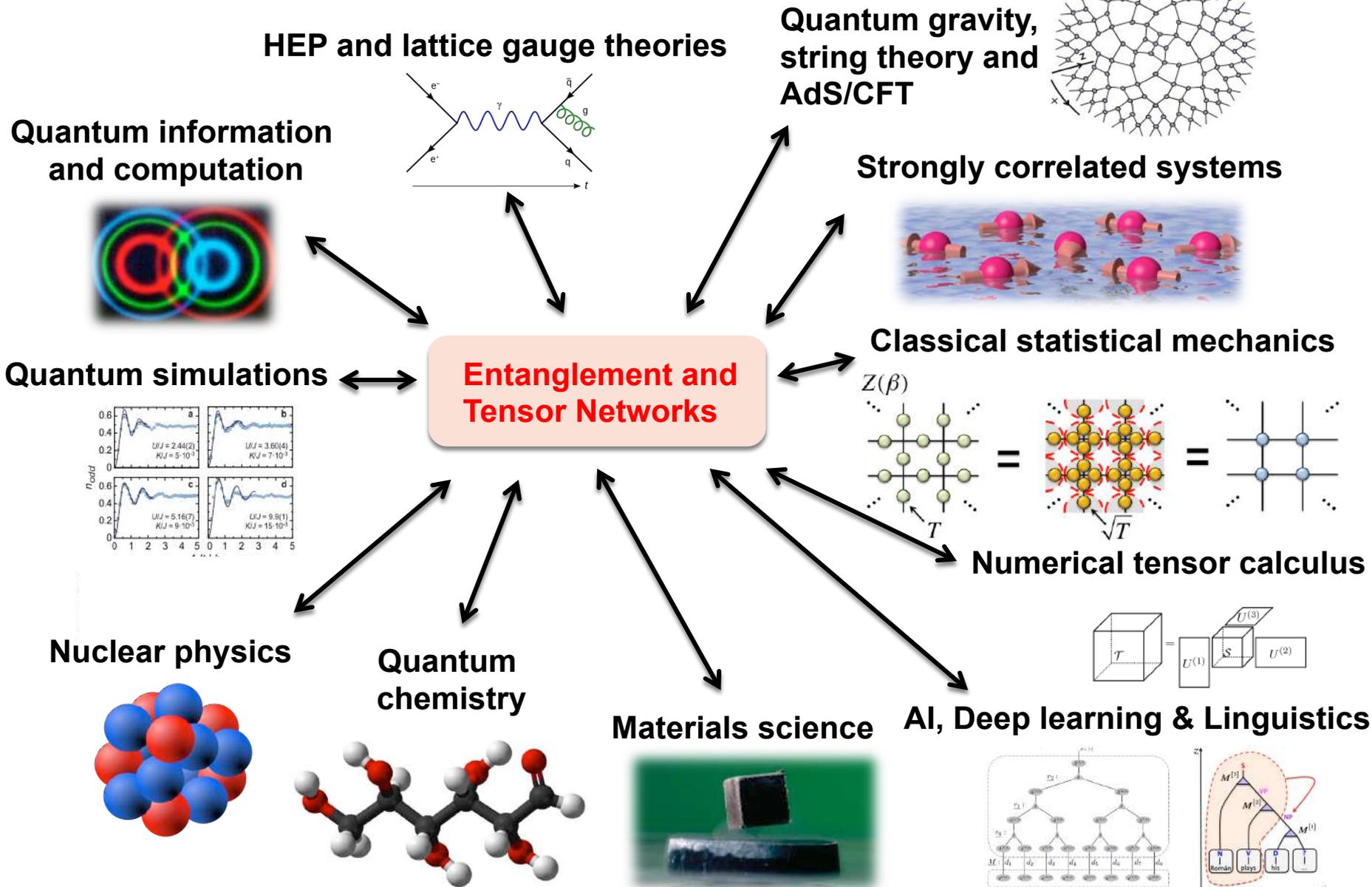
Efficient  $O(\text{poly}(N))$ , satisfy area-law, low-energy eigenstates of local Hamiltonians

# Tensor networks everywhere

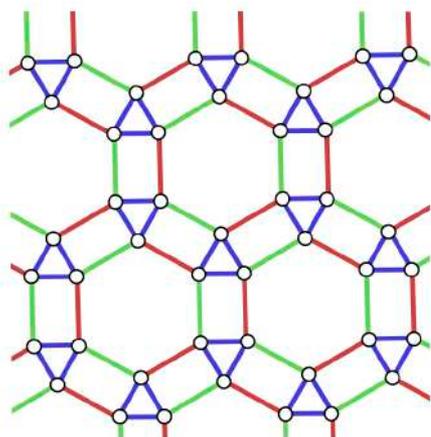


**Entanglement and  
Tensor Networks**

# Tensor networks everywhere



# iPEPS for the ruby model (1)



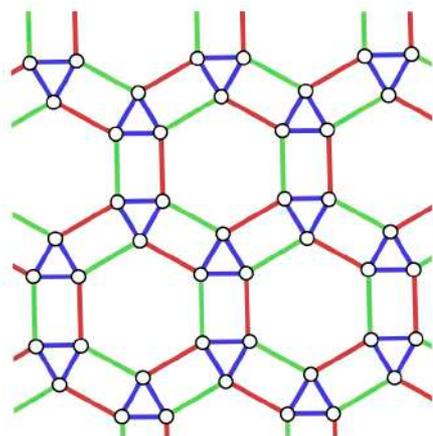
●—● YY

●—● XX

●—● ZZ

$$H = - \sum_{\alpha=x,y,z} J_{\alpha} \sum_{\alpha\text{-links}} \sigma_i^{\alpha} \sigma_j^{\alpha}$$

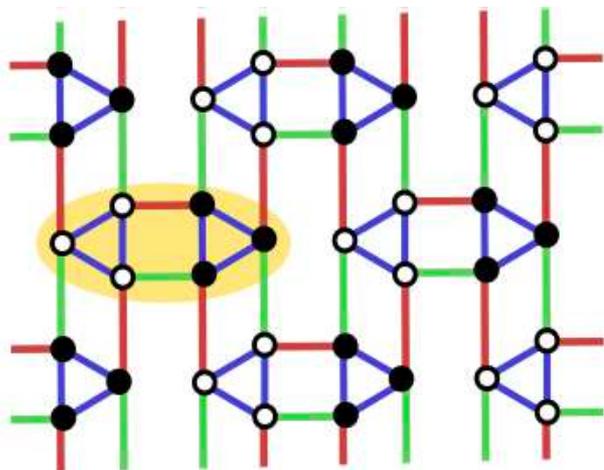
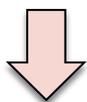
# iPEPS for the ruby model (1)



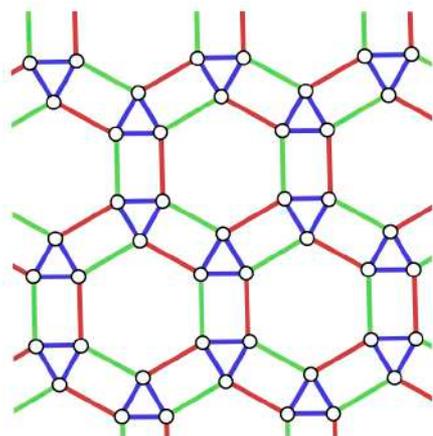
- YY
- XX
- ZZ

$$H = - \sum_{\alpha=x,y,z} J_{\alpha} \sum_{\alpha\text{-links}} \sigma_i^{\alpha} \sigma_j^{\alpha}$$

stretch



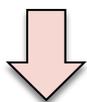
# iPEPS for the ruby model (1)



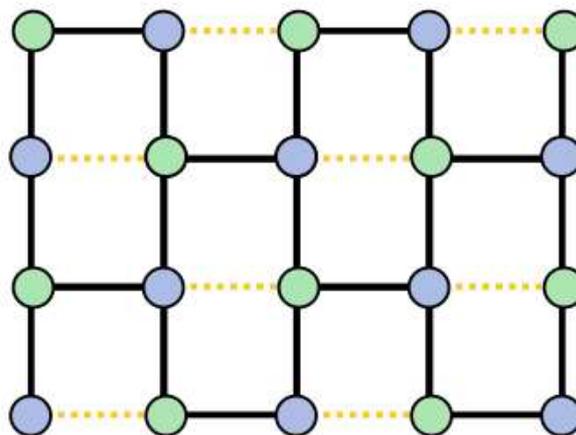
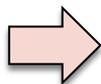
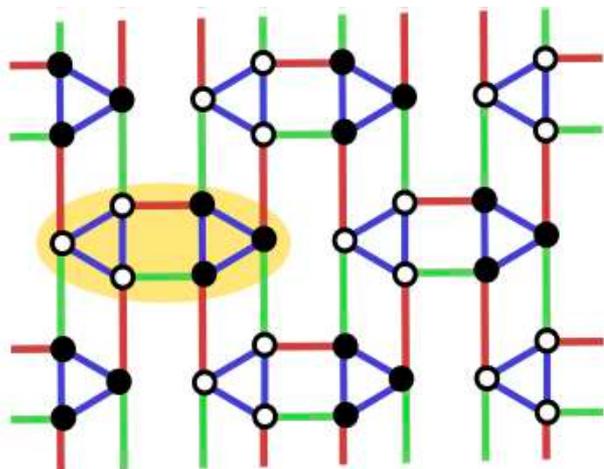
- YY
- XX
- ZZ

$$H = - \sum_{\alpha=x,y,z} J_{\alpha} \sum_{\alpha\text{-links}} \sigma_i^{\alpha} \sigma_j^{\alpha}$$

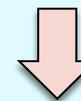
stretch



coarse-grain

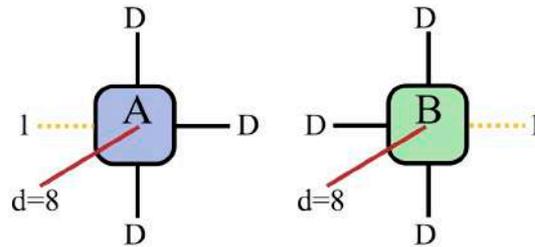
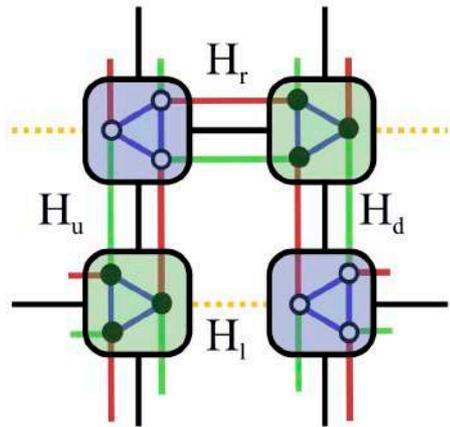


Coarse-graining triangles leads to brickwall lattice



Study with iPEPS

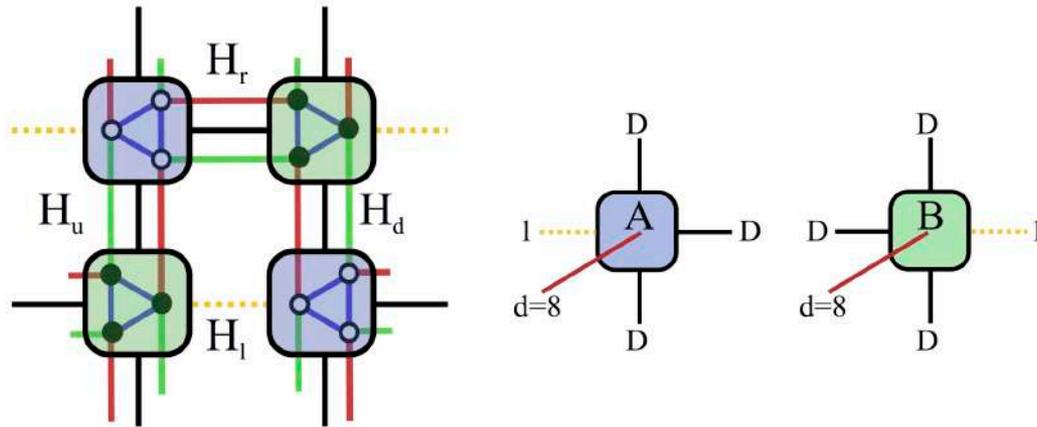
# iPEPS for the ruby model (2)



3-site coarse-graining

2 x 1 unit cell

# iPEPS for the ruby model (2)



3-site coarse-graining

2 x 1 unit cell

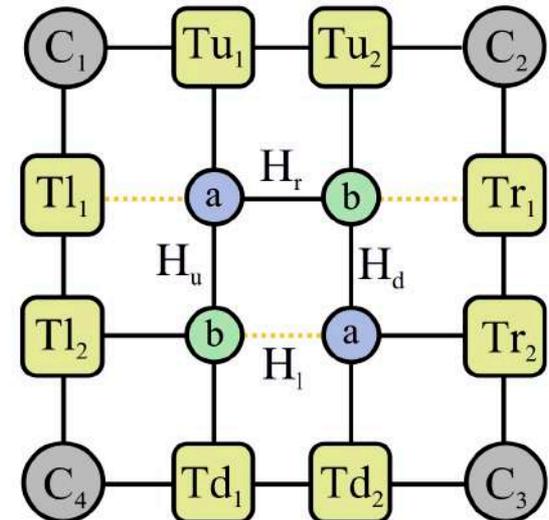
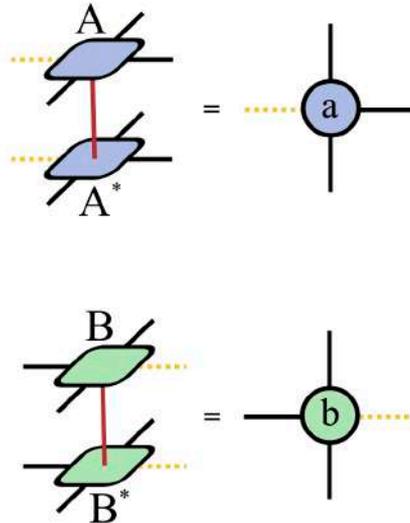
Effective environments  
computed with Corner  
Transfer Matrices (CTM)

*RO, G. Vidal, PRB 80 094403 (2009)*

Full update + gauge fixing

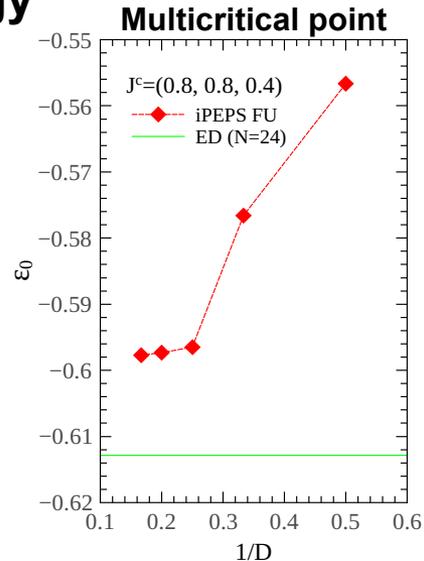
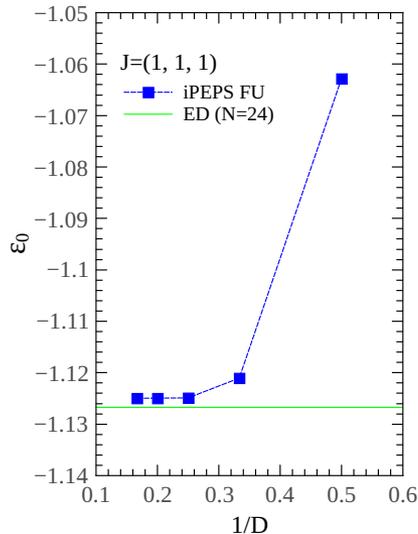
*H. N. Phien et al, PRB 92, 035142 (2015)*

$$D = 12, \chi = 100, d = 8$$



# Energies and phase diagram

## Ground state energy



Comparison to exact diag.  
for 24 sites.

Good convergence.  
Multicritical point challenging.

3 phases separated by **2nd order QPTs**.

QPTs meet at a multicritical point.

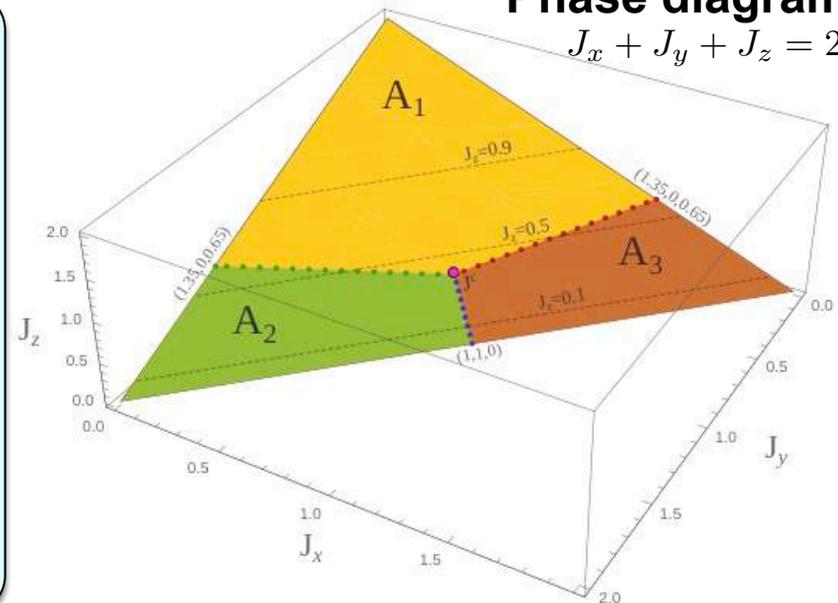
$A_1$  phase: gapped,  $Z_2 \times Z_2$  TO.

$A_2$  &  $A_3$  phases: gapless. Gapped by a perturbation, colored Ising anyons.

*S. S. Jahromi, M. Kargarian, S. F. Masoudi, A. Langari, PRB 94, 125145 (2016)*

## Phase diagram

$$J_x + J_y + J_z = 2$$



# Fidelity from CTMs (1)



## Ground state fidelity per site

*H.-Q. Zhou, RO, G. Vidal, PRL 100, 080601 (2008)*

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$

$$\ln d(\lambda_1, \lambda_2) \equiv \lim_{N \rightarrow \infty} \frac{\ln F(\lambda_1, \lambda_2)}{N}$$

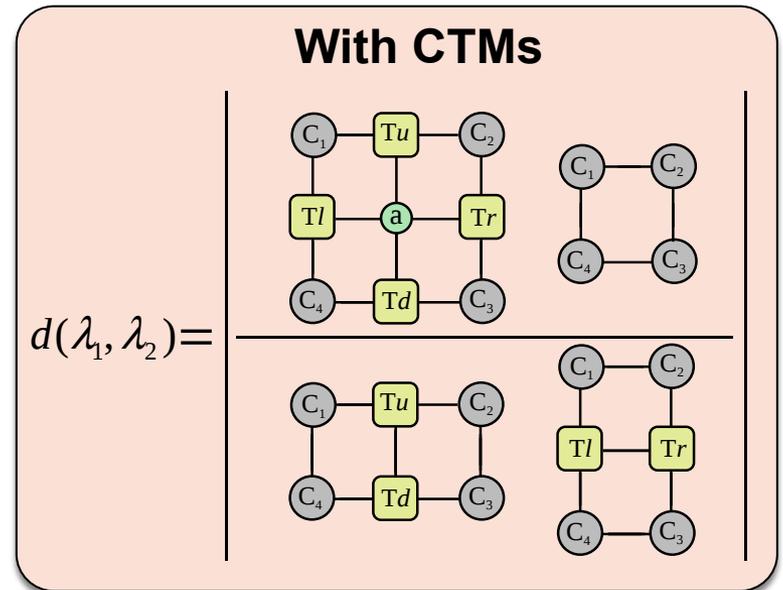
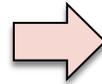
# Fidelity from CTMs (1)

## Ground state fidelity per site

*H.-Q. Zhou, RO, G. Vidal, PRL 100, 080601 (2008)*

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$

$$\ln d(\lambda_1, \lambda_2) \equiv \lim_{N \rightarrow \infty} \frac{\ln F(\lambda_1, \lambda_2)}{N}$$



(I'll prove it later)

# Fidelity from CTMs (1)

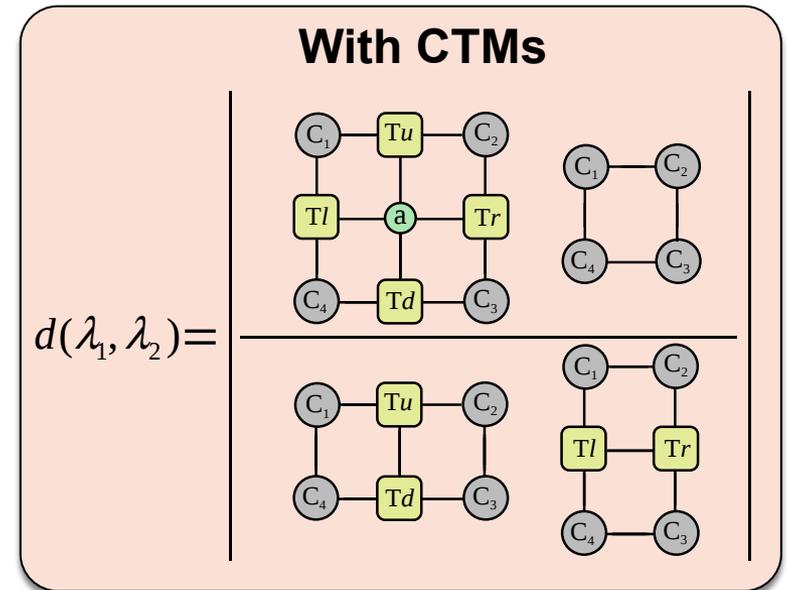
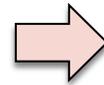
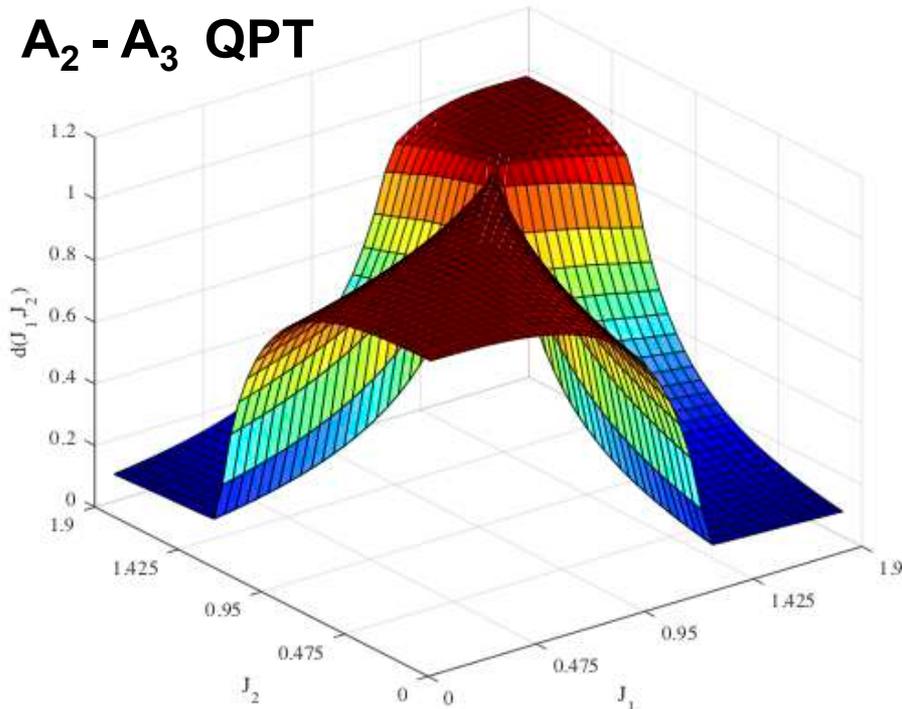
## Ground state fidelity per site

*H.-Q. Zhou, RO, G. Vidal, PRL 100, 080601 (2008)*

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$

$$\ln d(\lambda_1, \lambda_2) \equiv \lim_{N \rightarrow \infty} \frac{\ln F(\lambda_1, \lambda_2)}{N}$$

## A<sub>2</sub> - A<sub>3</sub> QPT



(I'll prove it later)

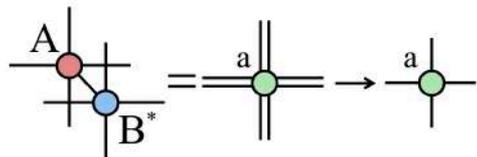
- Compatible with other observables:
- Entanglement entropy
  - 2-point correlators

# Fidelity from CTMs (proof of formula)

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$

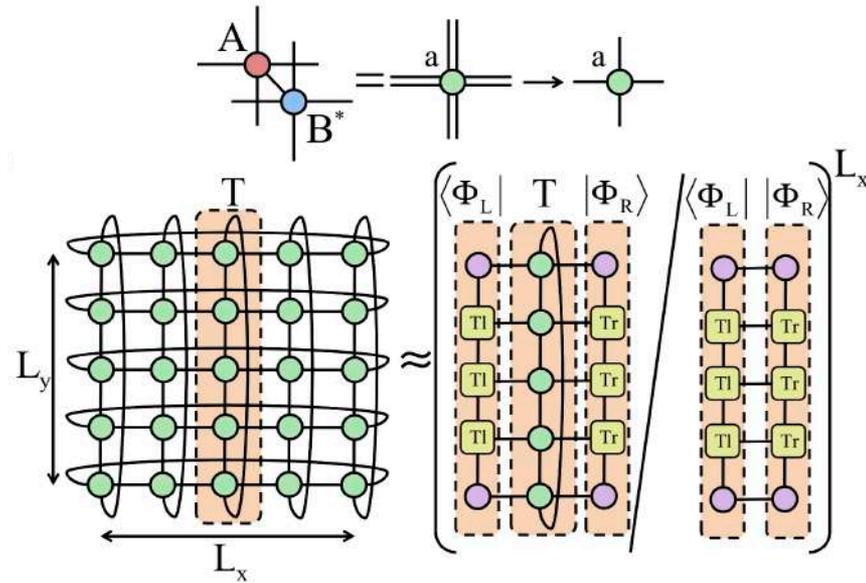
# Fidelity from CTMs (proof of formula)

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$



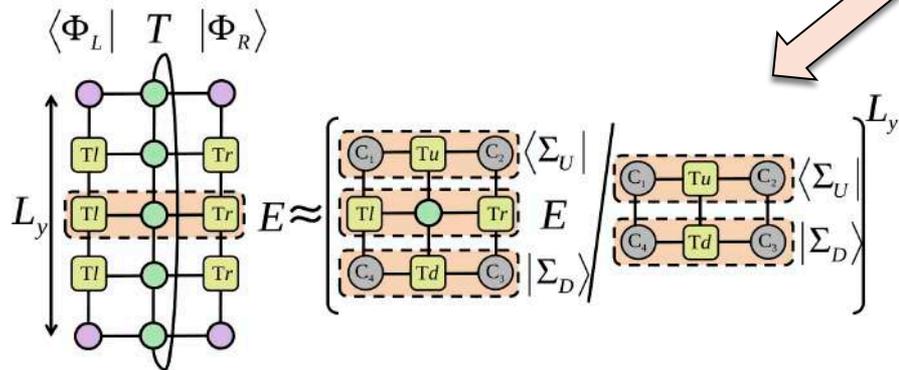
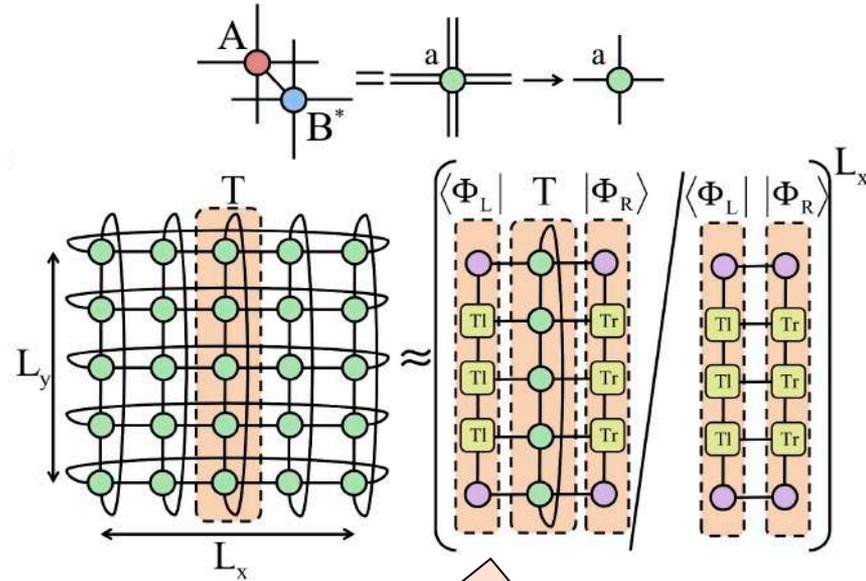
# Fidelity from CTMs (proof of formula)

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$



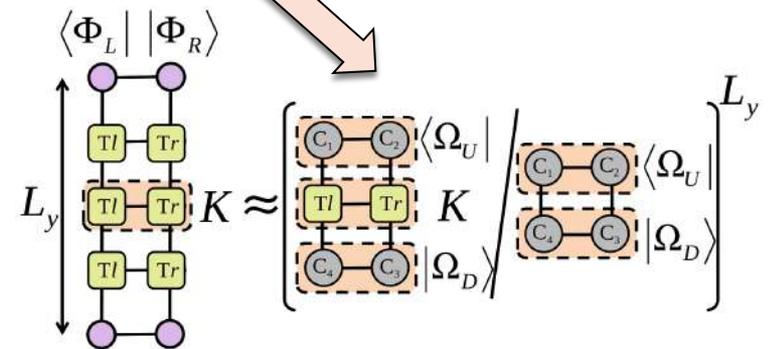
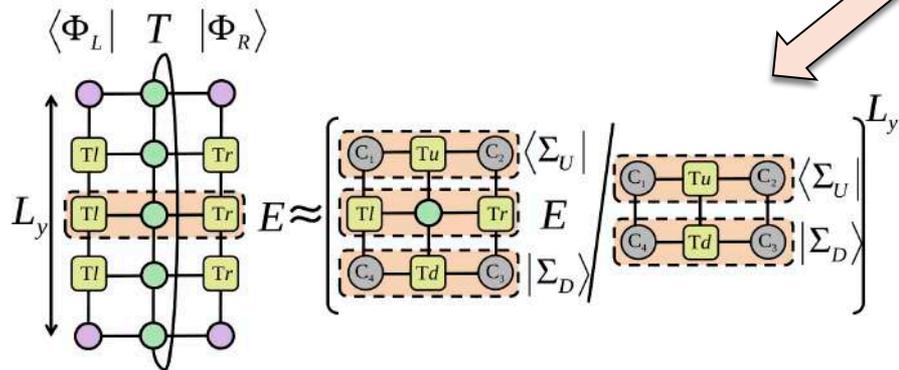
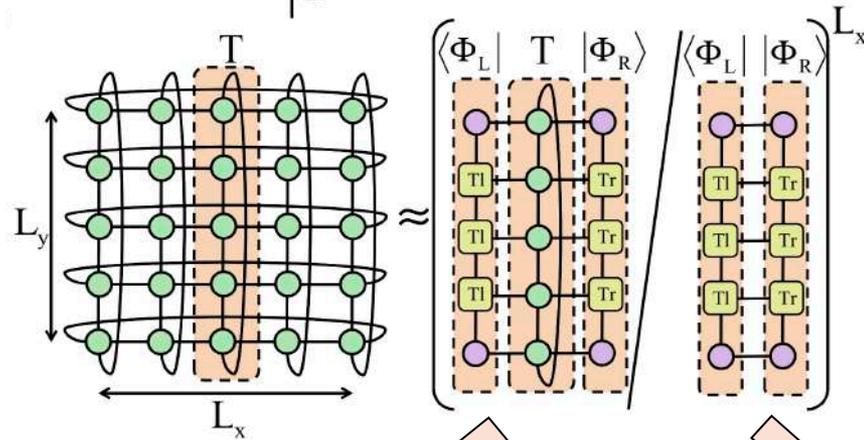
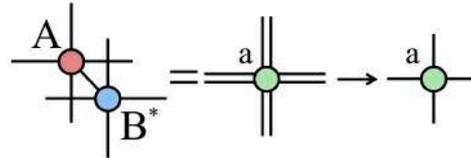
# Fidelity from CTMs (proof of formula)

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$



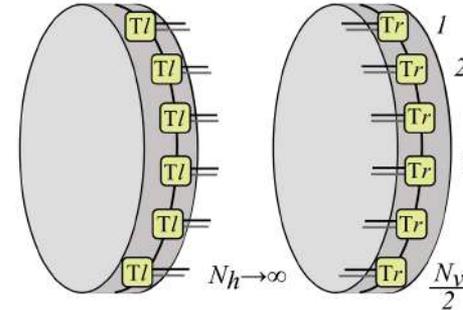
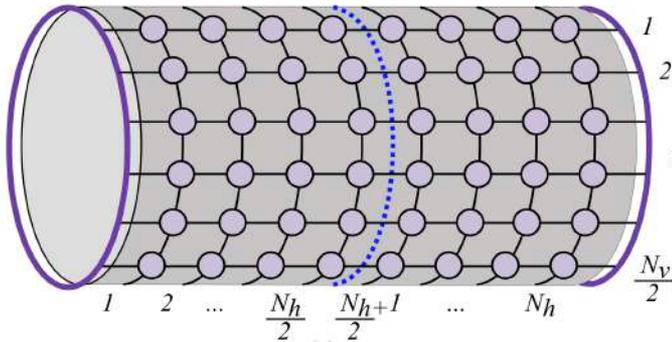
# Fidelity from CTMs (proof of formula)

$$F(\lambda_1, \lambda_2) = |\langle \Psi(\lambda_2) | \Psi(\lambda_1) \rangle| \sim d(\lambda_1, \lambda_2)^N$$

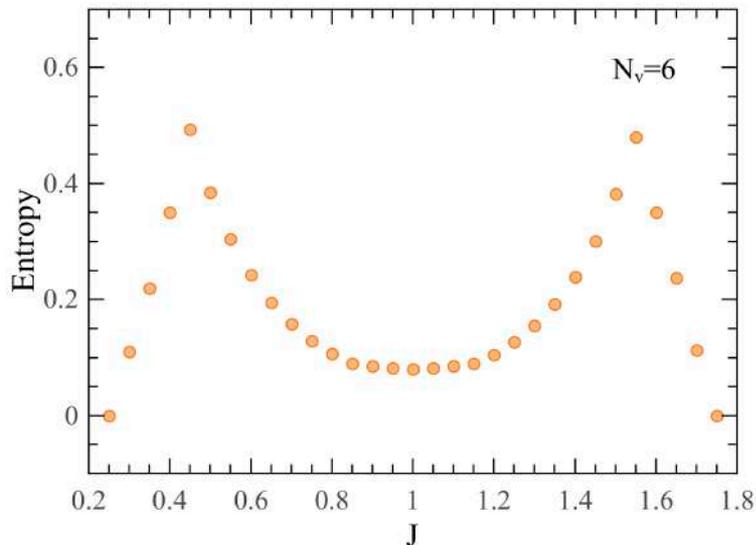


# Entanglement entropy on cylinders

J. I. Cirac et al, PRB 83, 245134 (2011)



$$\rho = U \sqrt{\sigma_L^T} \sigma_R \sqrt{\sigma_L^T} U^\dagger \quad (\text{for half an infinite cylinder})$$



$A_2 - A_1 - A_3$  transitions

Scaling with  $N_v$ : non-zero topological entropy, but hard to identify

Edge Hamiltonian also hard to identify

# OUTLOOK

- 1) Phase diagram and phase transitions of the topological ruby model with iPEPS
- 2) Ground state energies, local fidelities, entanglement entropy of half-infinite cylinders
- 3) Open: detailed numerical characterization of topological orders (specially for the non-abelian phases), and the multicritical point. Phase diagram in magnetic field?
- 4) Open: study via combined approaches? (e.g., pCUT + iPEPS)

*S. Dusuel et al, PRL 106, 107203 (2011)*

# Other recent activities at our group

- Quantum antiferromagnets on the star lattice

*S. S. Jahromi, RO, PRB 98 155108 (2018)*

- gPEPS algorithm for any lattice

*S. S. Jahromi, RO, arXiv:1808.00680*

- SU(2) tensor network algorithms

*P. Schmoll et al, arXiv:1809.08180*

- 2d annealing with tensor networks

*A. Kshetrimayum et al, arXiv:1809.08258*

- Quantum criticality on a chiral ladder

*P. Schmoll et al, arXiv:1812.01311*

- Topological order on the Bloch sphere

*R. Liss, T. Mor, RO, arXiv:1812.00671*

- Quantum computing for finance

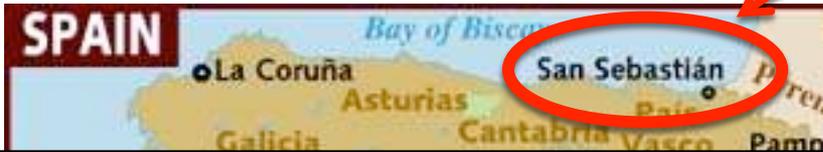
*RO, S. Mugel, E. Lizaso, arXiv: 1807.03890, arXiv:1810.07690*

















San Sebastián



Applications welcomed

