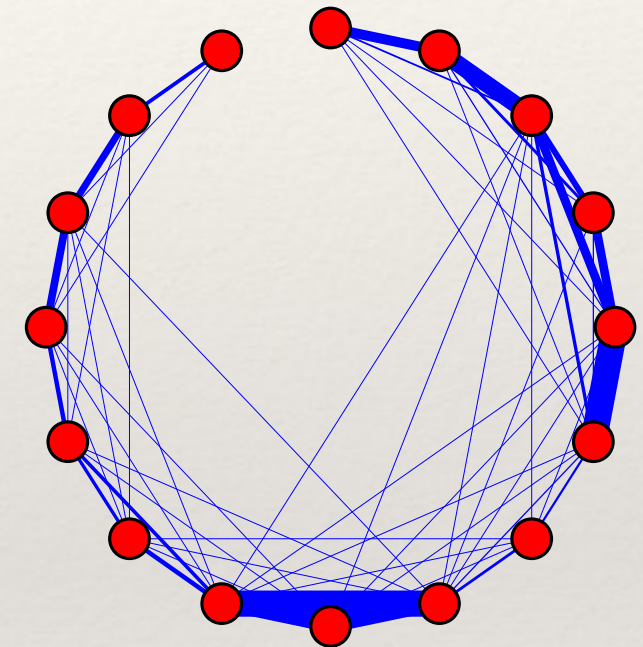
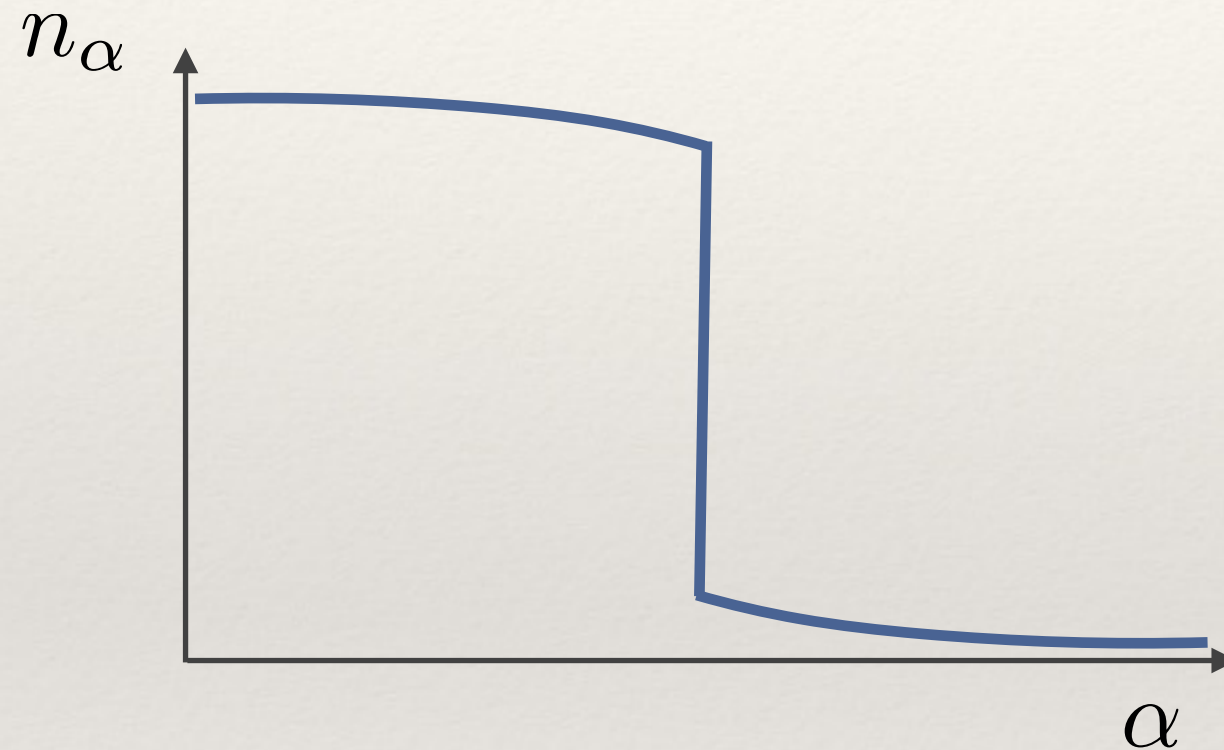


Many-body localization characterized by entanglement and occupations of natural orbitals

Jens Hjörleifur Bärđarson

Max Planck Institute PKS, Dresden

KTH Royal Institute of Technology, Stockholm





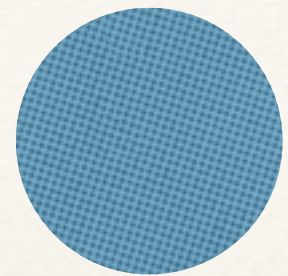
I: Broken symmetry



II: Adiabatic continuity & renormalization

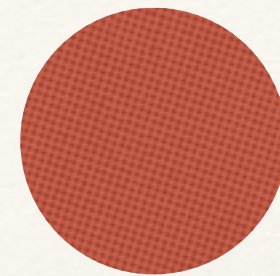


Fermi liquid theory

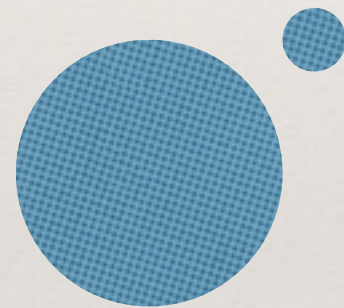


Fermi Gas

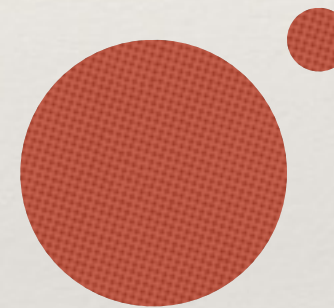
Ground state



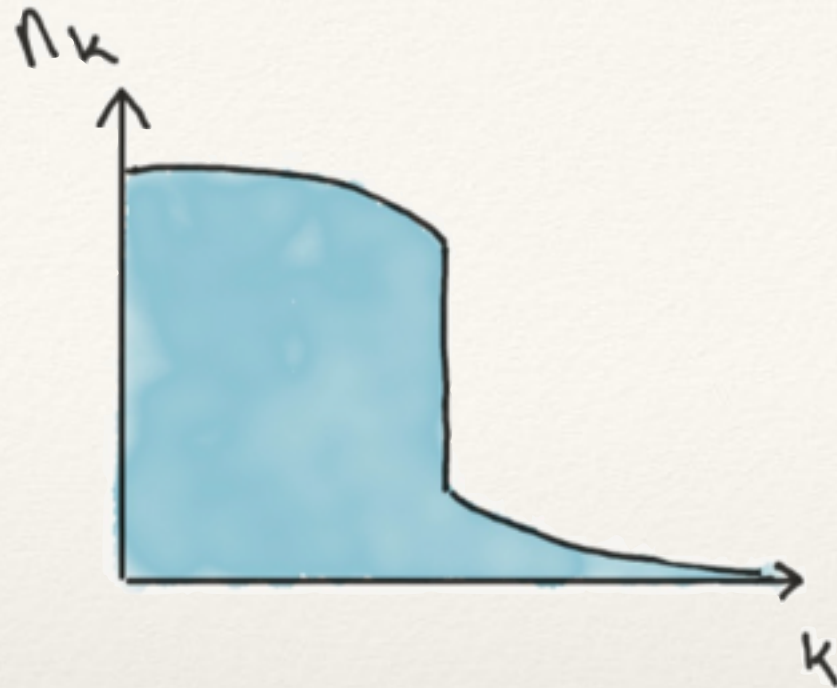
Fermi Liquid



Quasiparticles



Fermi liquid theory

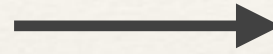


$$\mathcal{E} = \mathcal{E}_0 + \sum_{\mathbf{k}\sigma} (E_{\mathbf{k}\sigma}^{(0)} - \mu) \delta n_{\mathbf{k}\sigma} + \frac{1}{2} \sum_{\mathbf{k}\mathbf{k}', \sigma\sigma'} f_{\mathbf{k}\sigma, \mathbf{k}'\sigma'} \delta n_{\mathbf{k}\sigma} \delta n_{\mathbf{k}'\sigma'} + \dots$$

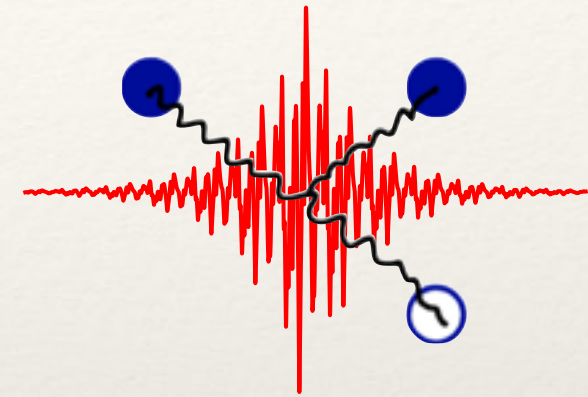
$$c_{\mathbf{k}\sigma}^\dagger = \sqrt{Z_{\mathbf{k}}} a_{\mathbf{k}\sigma}^\dagger + \sum A(\mathbf{k}_4\sigma_4, \mathbf{k}_3\sigma_3; \mathbf{k}_2\sigma_2, \mathbf{k}\sigma) a_{\mathbf{k}_4\sigma_4}^\dagger a_{\mathbf{k}_3\sigma_3}^\dagger a_{\mathbf{k}_2\sigma_2} + \dots$$

Many-body localization

Anderson Insulator



Many-body localized insulator



Basko, Aleiner, Altshuler, Ann. Phys 2006

see also: Altshuler, Gefen, Kamenev, Levitov PRL 1997 Gornyi, Mirlin, Polyakov PRL 2005

Anderson insulator (many-body)

$$H_0 = -\frac{1}{2} \sum_{i=1}^L \left(c_{i+1}^\dagger c_i + c_i^\dagger c_{i+1} + \epsilon_i n_i \right) = \sum_i \epsilon_i a_i^\dagger a_i \quad \epsilon_i \in [-W, W]$$



$$\tilde{n}_i = a_i^\dagger a_i$$

$$[\tilde{n}_i, H_0] = 0$$

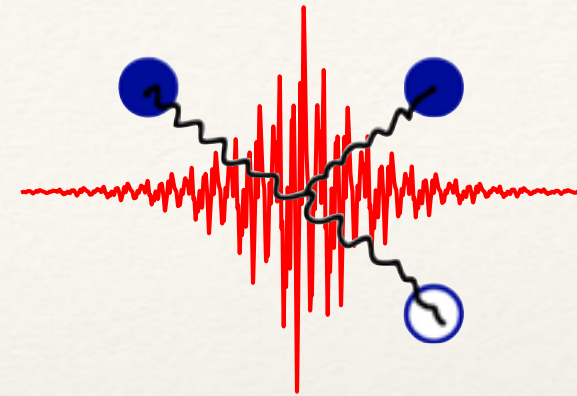
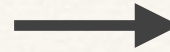
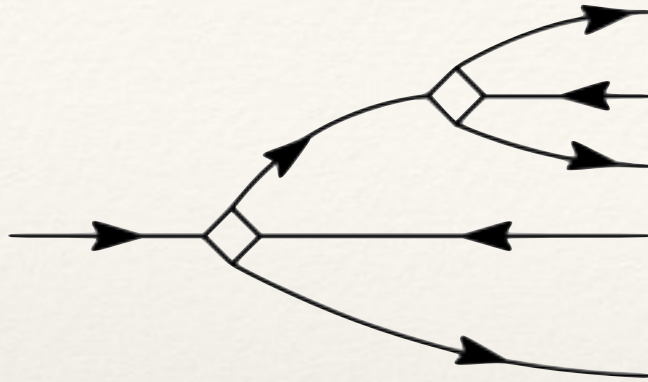
Local integrals of motion

$$|\tilde{n}_1, \tilde{n}_2, \dots, \tilde{n}_L\rangle$$

Eigenstates are product states

Emergent integrability in many-body localization

$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} a_i^\dagger a_j^\dagger a_l a_k$$



$$H = \sum_i h_i n_i + \sum_{ij} J_{ij} n_i n_j + \dots$$

where

$$J_{ij} \sim e^{-|x_i - x_j|/\xi^*}$$

$$[n_i, H] = 0$$

Local conserved integrals of motion

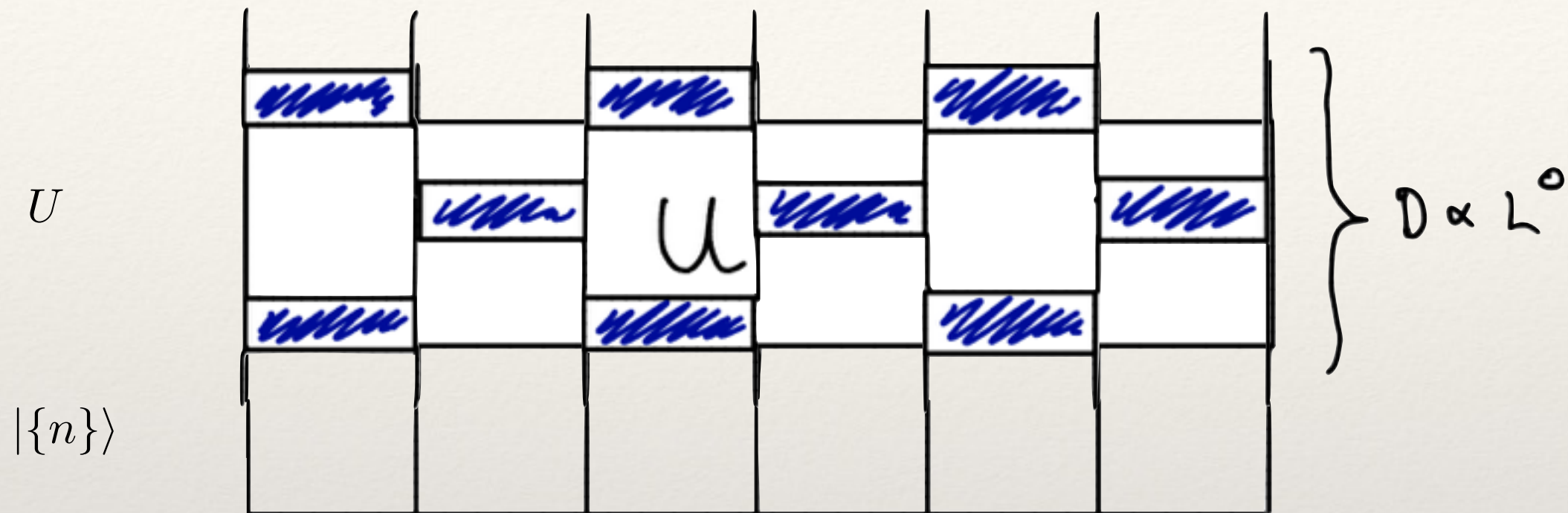
$$|n_1, n_2, \dots, n_L\rangle$$

All eigenstates product states

$$c_i = \sum_j Z_j^i a_j + \sum_{jkl} f_{jkl}^i a_j^\dagger a_k a_l + \dots$$

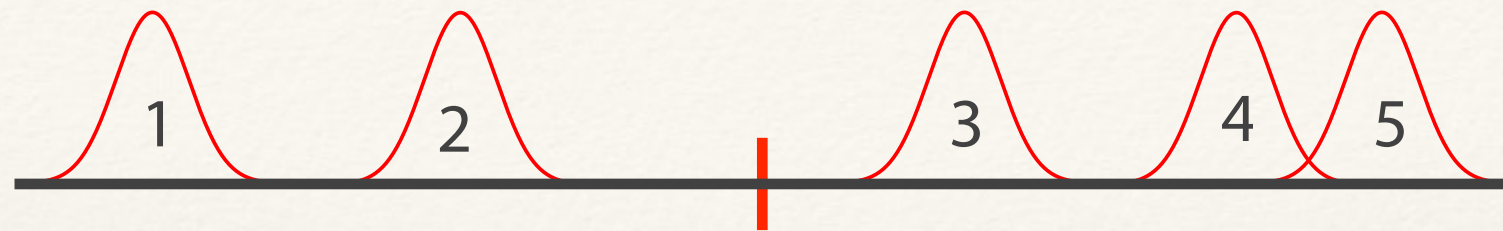
Quasiparticles

Every MBL eigenstate related to an Anderson insulator eigenstate via a finite depth quantum circuit



Many-body localized eigenstates have an area law entanglement

$$S(A) = -\text{Tr}_A \rho_A \log \rho_A \quad \rho_A = \text{Tr}_{\bar{A}} |\psi\rangle\langle\psi|$$

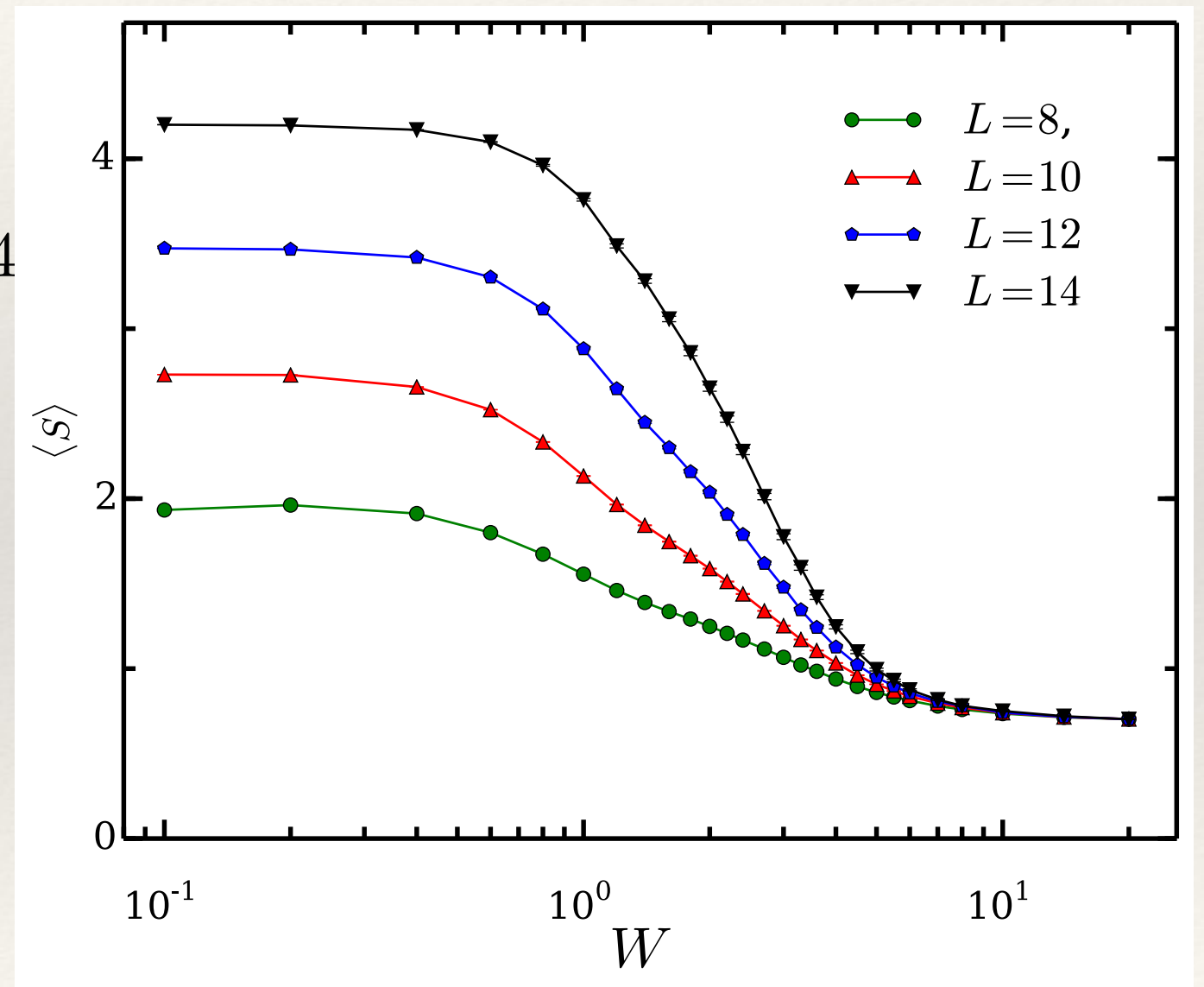


$$|\psi\rangle = c_1^\dagger c_2^\dagger c_3^\dagger c_4^\dagger c_5^\dagger |0\rangle = |12345\rangle$$

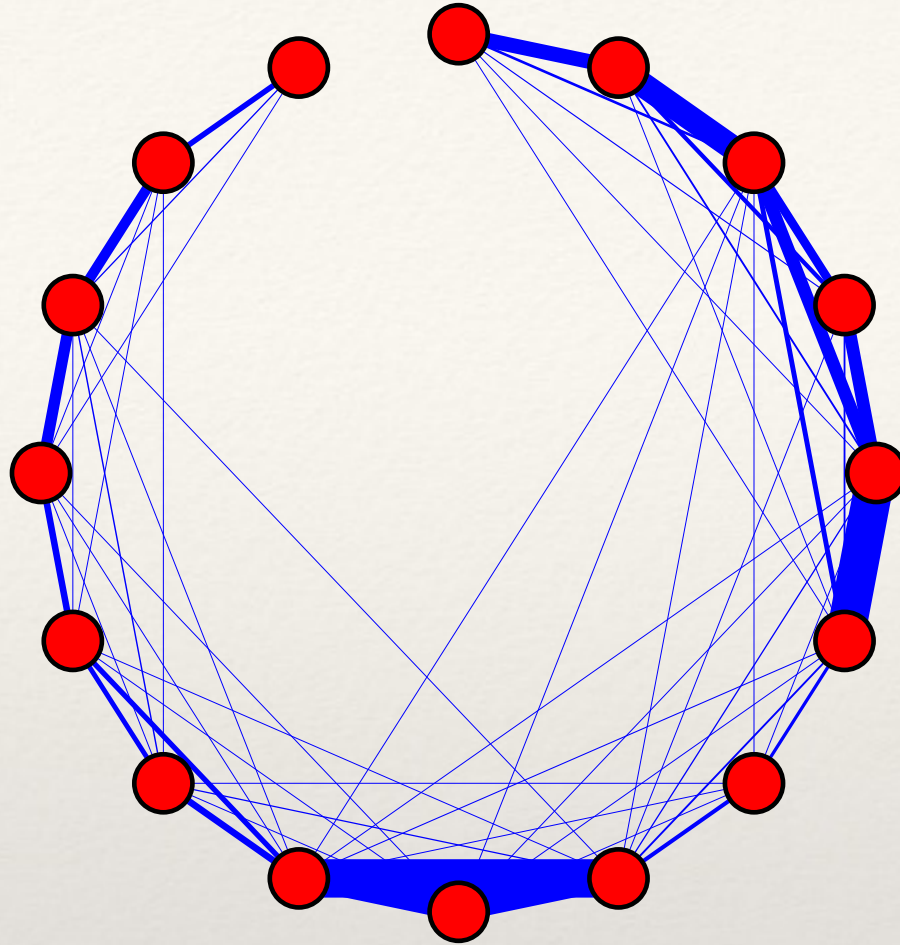
$$\rho_R = \frac{1}{L} \sum_{i \in R} |\psi\rangle\langle\psi| = \frac{1}{L} \sum_{i \in R} |i\rangle\langle i|$$

$$S_{\text{MBL}} \sim \xi$$

$$S_{\text{Page}} = \frac{L}{2} \log 2 - \frac{1}{2}$$



Mutual information directly reveals the entanglement structure of eigenstates



$$\mathcal{I}(A, B) = S(A) + S(B) - S(A \cup B)$$

One particle density matrix — definitions

$$H|\psi_n\rangle = E_n|\psi_n\rangle$$

$$\rho_{ij} = \langle\psi_n|c_i^\dagger c_j|\psi_n\rangle$$

$$\rho|\phi_\alpha\rangle = n_\alpha|\phi_\alpha\rangle$$

$|\phi_\alpha\rangle$

Natural orbitals

n_α

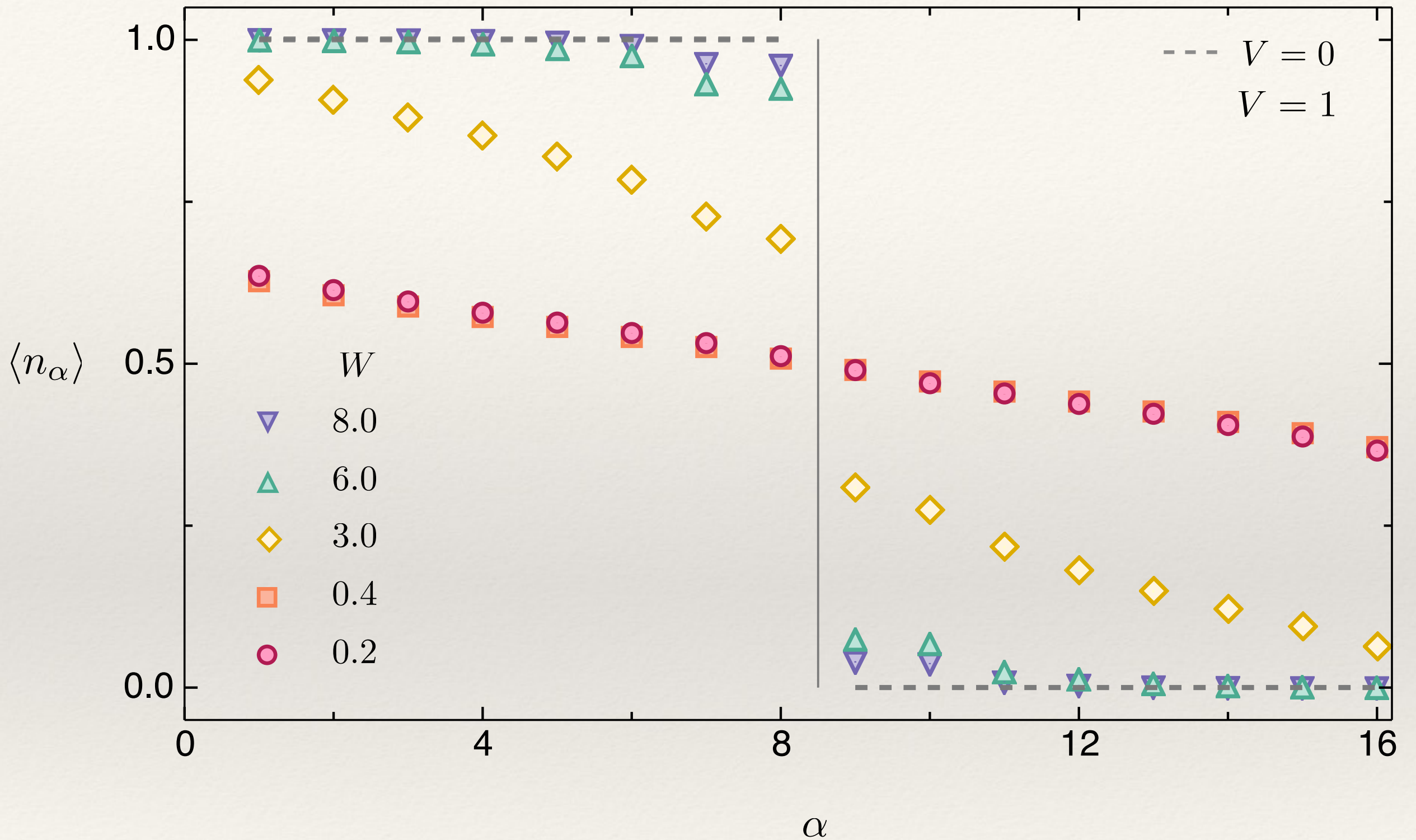
Occupations

$$n_1 \geq n_2 \geq \dots \geq n_L$$

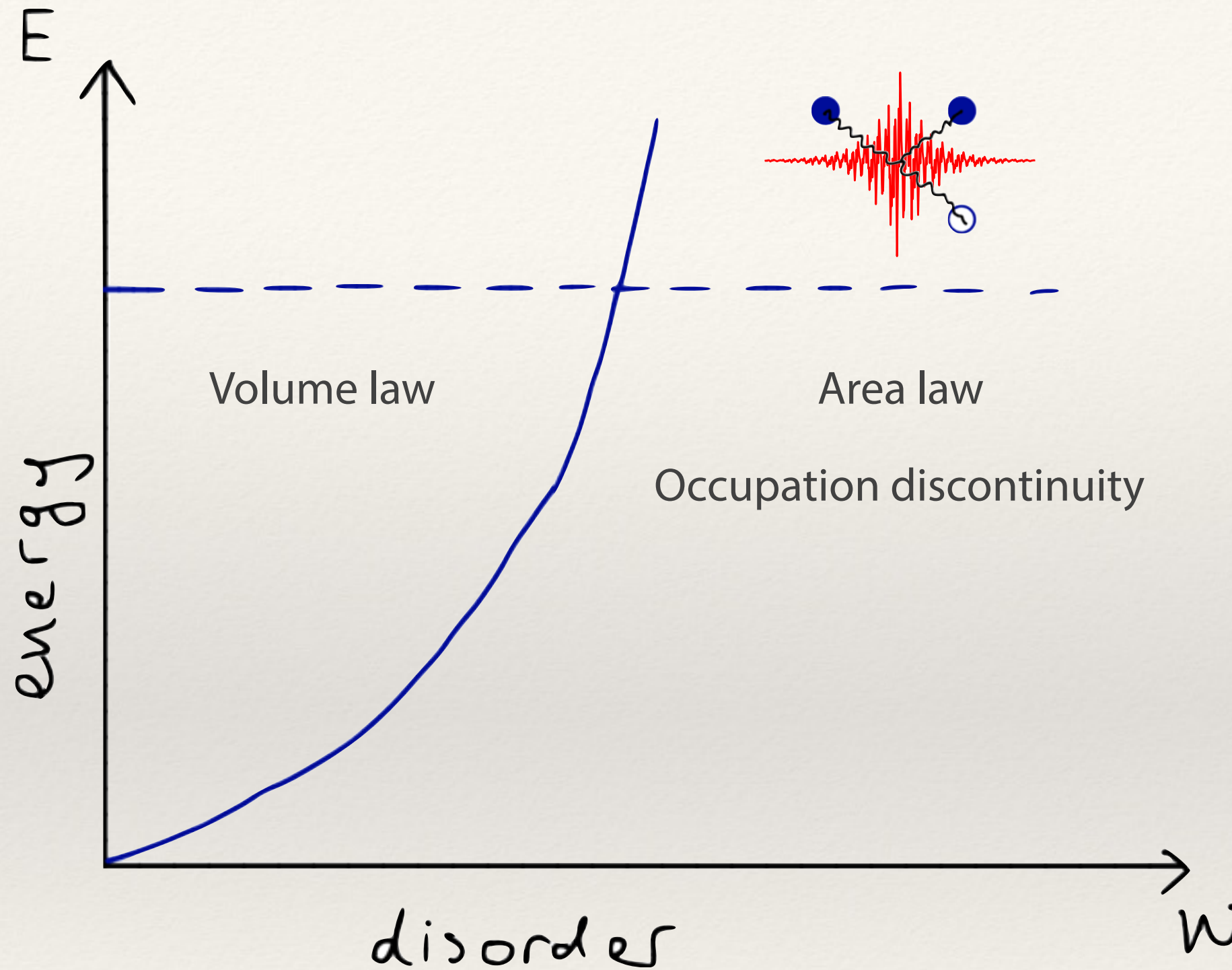
$$\text{tr } \rho = \sum_{\alpha=1}^L n_\alpha = N$$

Occupations

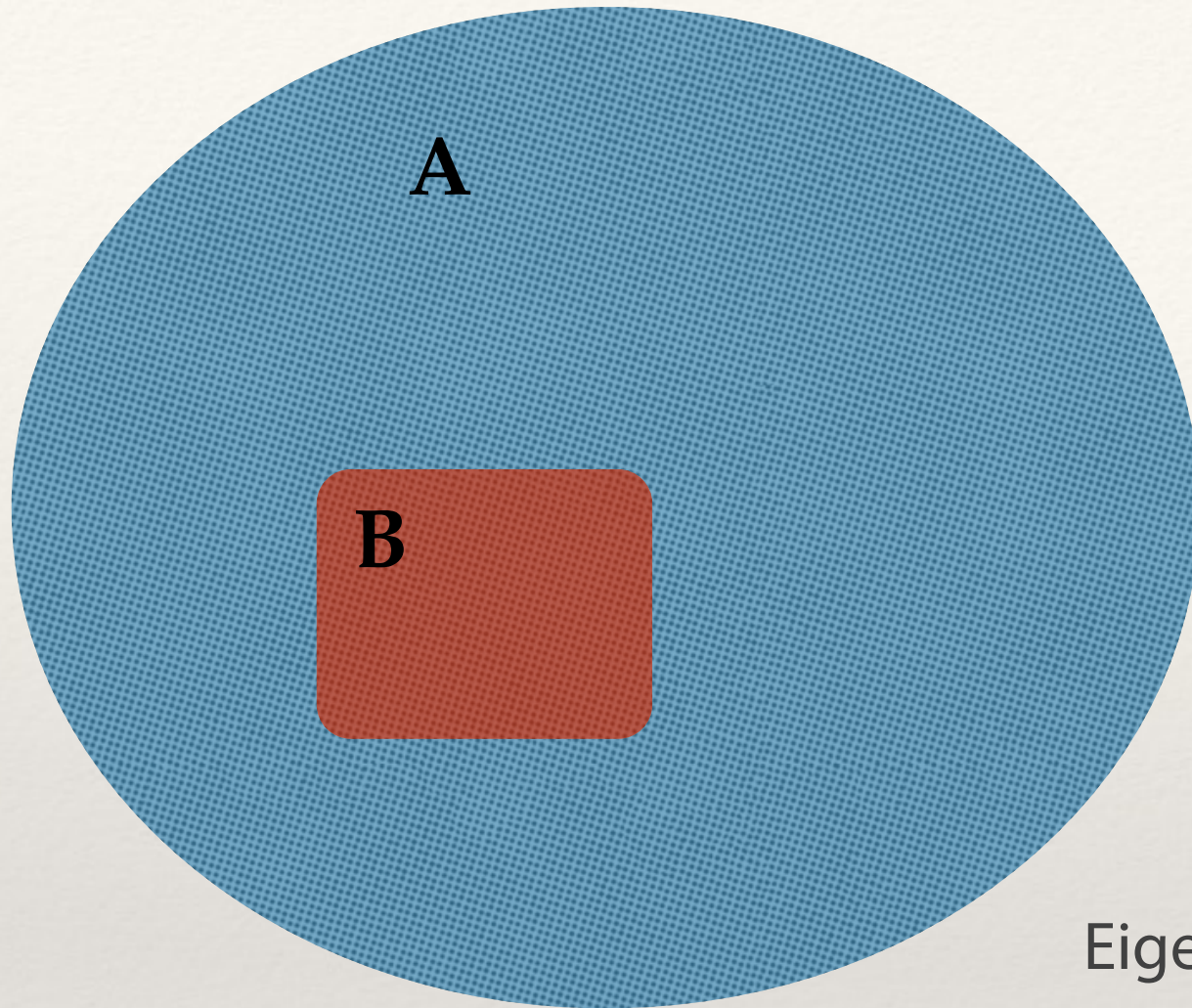
$$H = t \sum_{i=1}^L \left[-\frac{1}{2} (c_i^\dagger c_{i+1} + \text{h.c.}) + \epsilon_i \left(n_i - \frac{1}{2} \right) + V \left(n_i - \frac{1}{2} \right) \left(n_{i+1} - \frac{1}{2} \right) \right]$$



Many-body localization as eigenstate quantum phase transition



Entanglement entropy



$$|\psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$$

$$\rho_B = \text{Tr}_A |\psi\rangle\langle\psi|$$

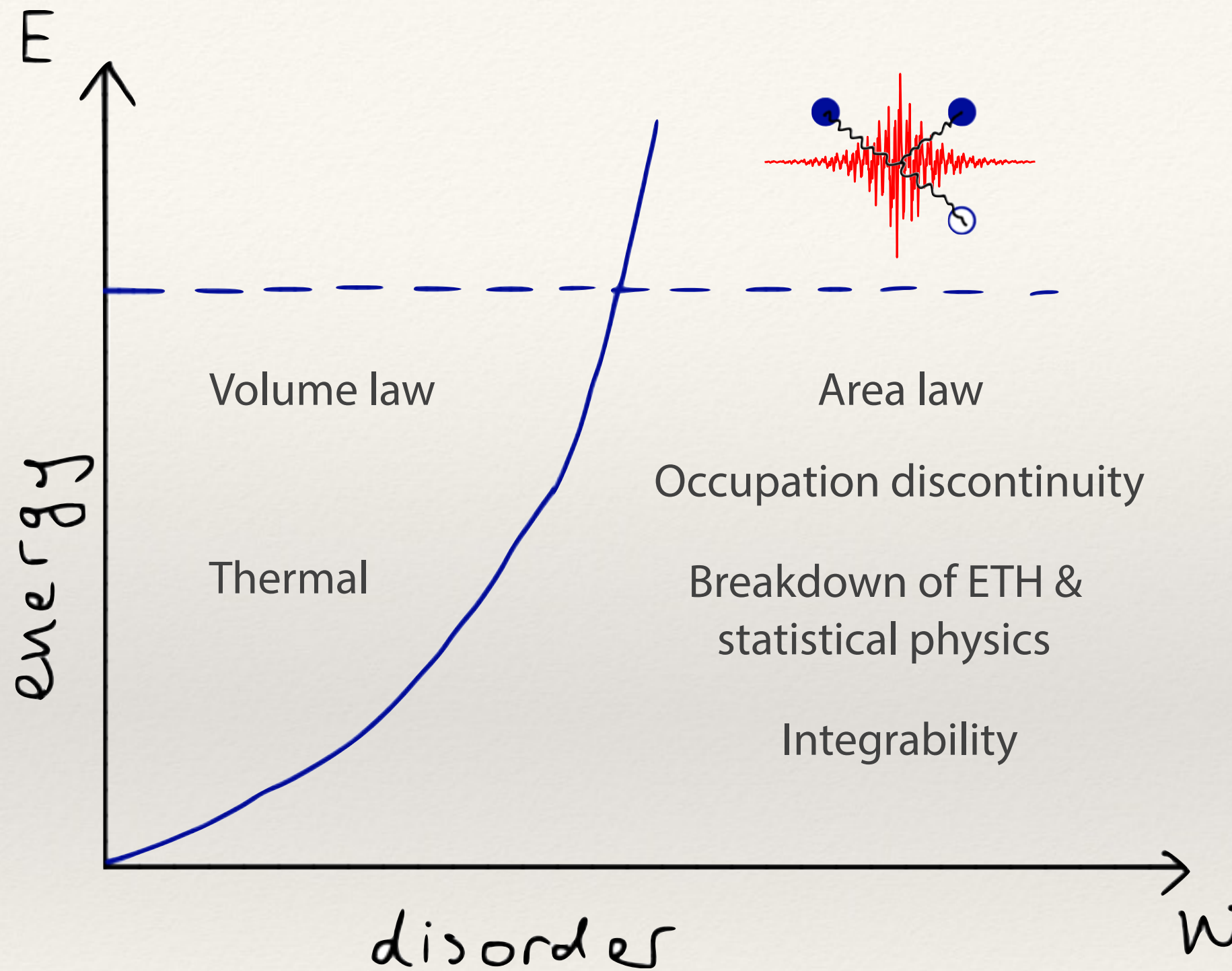
$$S = -\text{Tr} \rho_B \log \rho_B$$

Eigenstate thermalization hypothesis

$$H|\psi\rangle = E|\psi\rangle$$

$$\rho_B \sim e^{-\beta H_B}$$

Many-body localization as eigenstate quantum phase transition

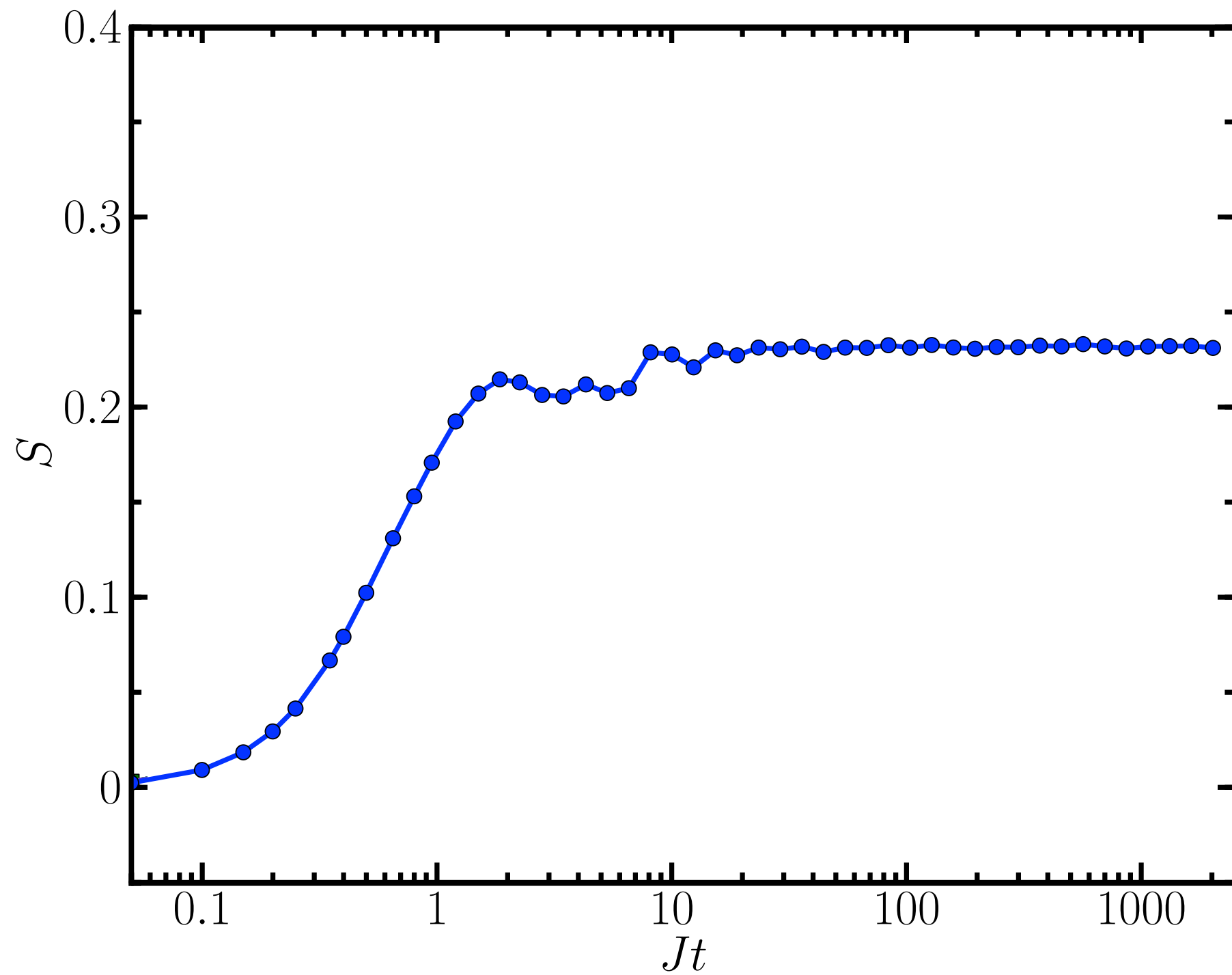


Global quench from a charge density wave — Anderson insulator

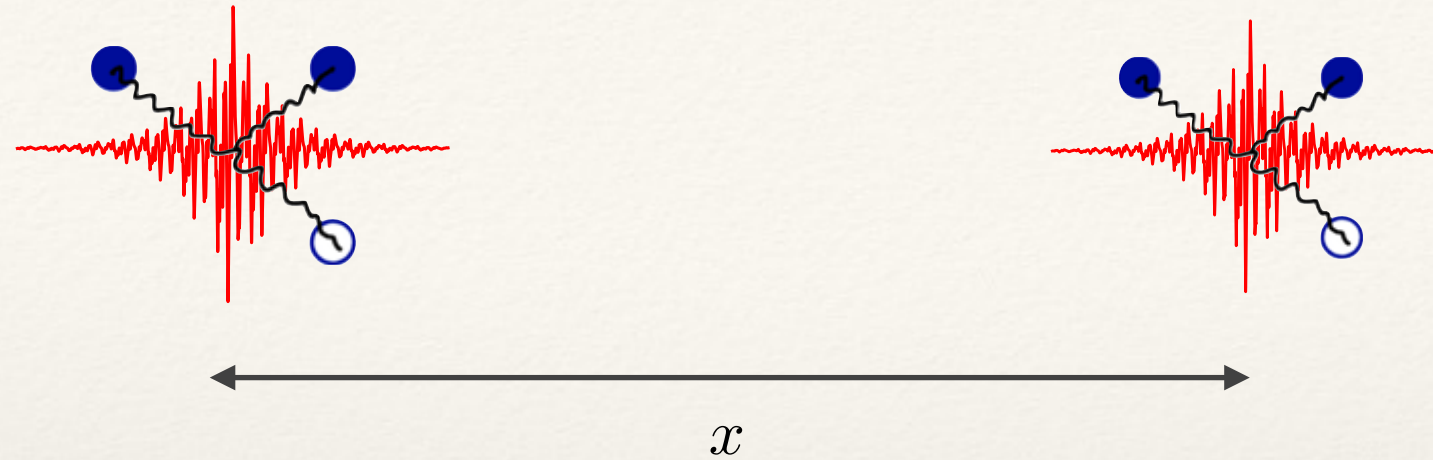


$$|\psi(t)\rangle = \exp(-iHt) \prod_{i=1}^{L/2} c_{2i}^\dagger |0\rangle$$

Entanglement vs. time — starting from a product state



Quasiparticle interaction leads to dephasing in dynamics



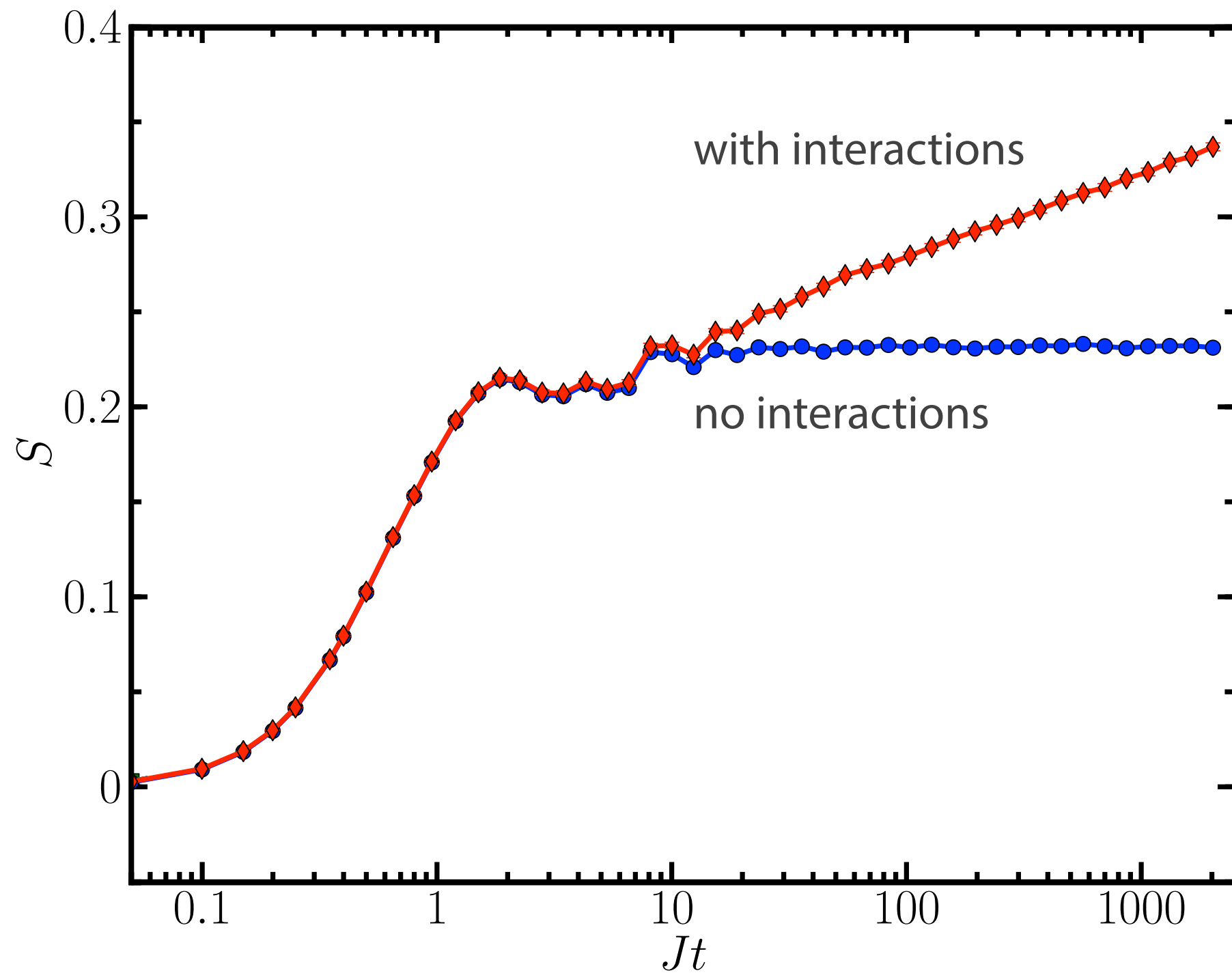
$$J_{ij} \sim J_0 e^{-|x_i - x_j|/\xi^*}$$

$$t_d \sim J_0^{-1} e^{x/\xi^*}$$

$$x_d(t) \sim \xi^* \log(J_0 t)$$

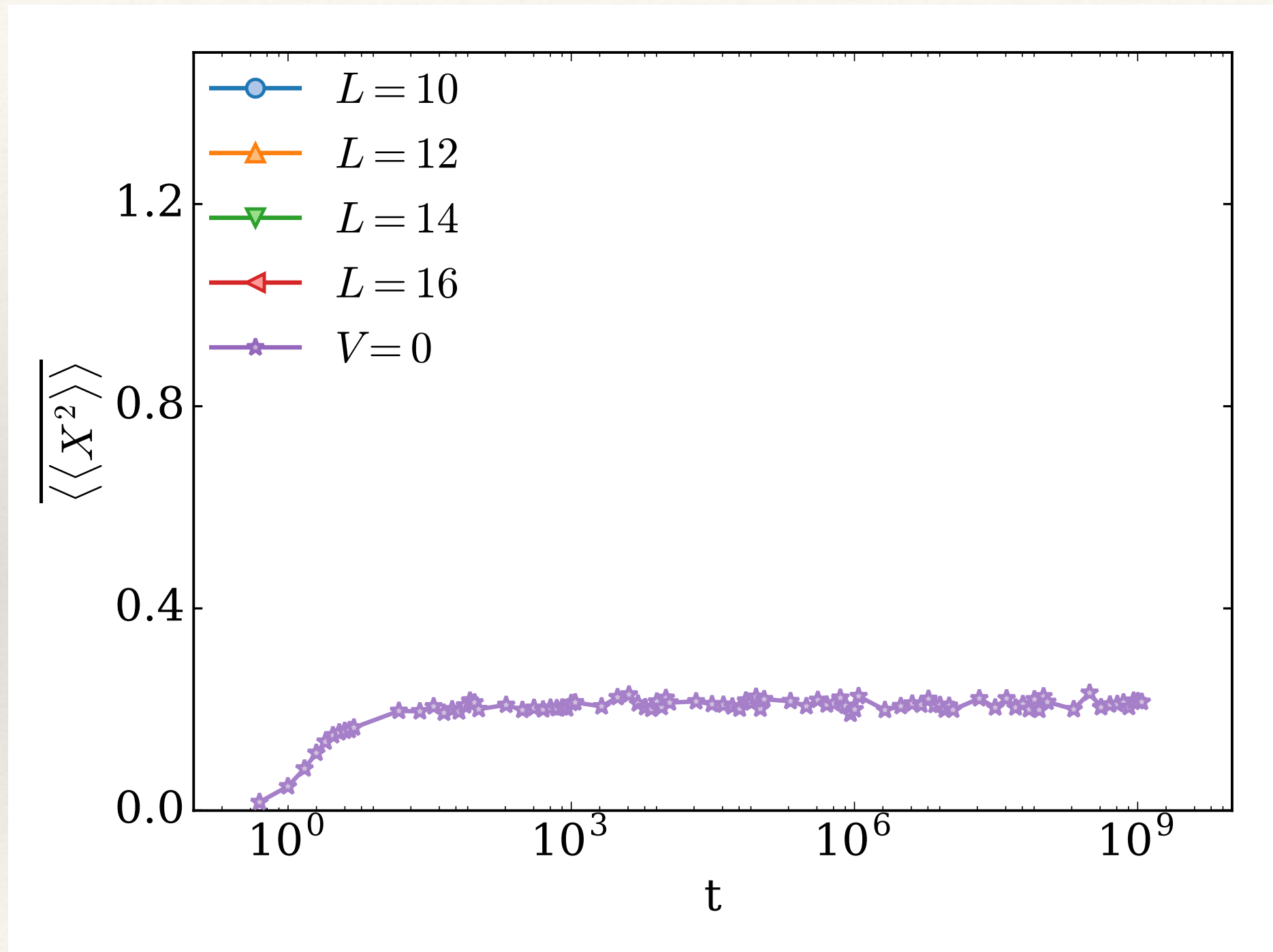
$$S(t) \sim \xi^* \log(J_0 t)$$

Entanglement vs. time — starting from a product state

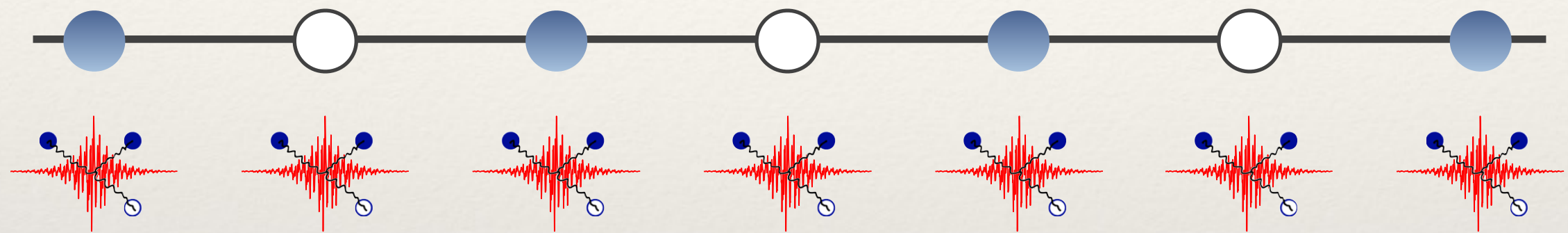


Mutual information after quench from product state

$$\langle\langle X^2 \rangle\rangle = \sum_j j^2 \mathcal{I}_j(t) - \left(\sum_j j \mathcal{I}_j \right)^2 \quad \mathcal{I}_j = \mathcal{I}(A = [1], B = [j])$$



Charge density wave provides partial occupation of quasiparticles

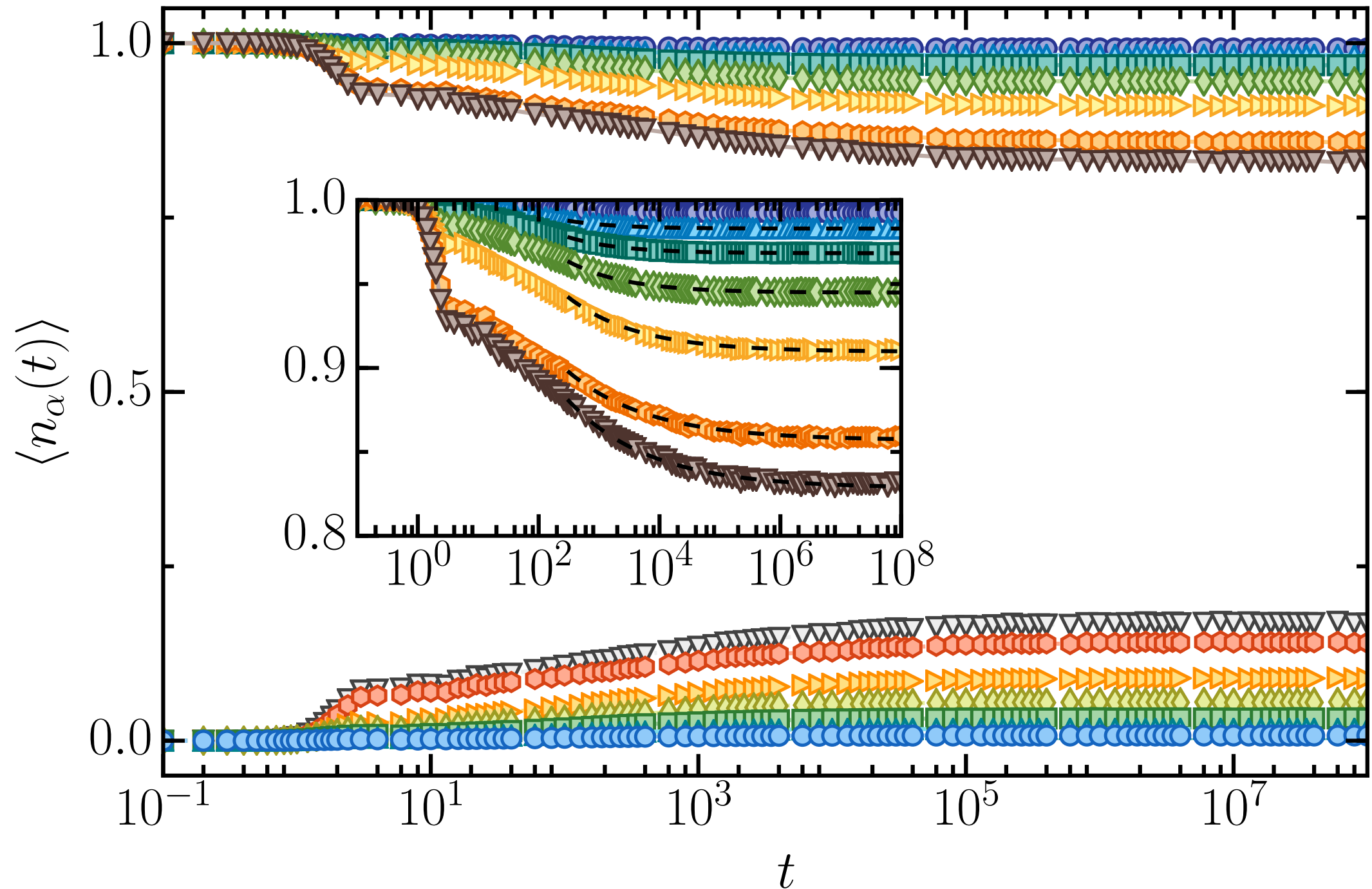


$$|\text{CDW}\rangle \sim |1010101\rangle + \dots$$

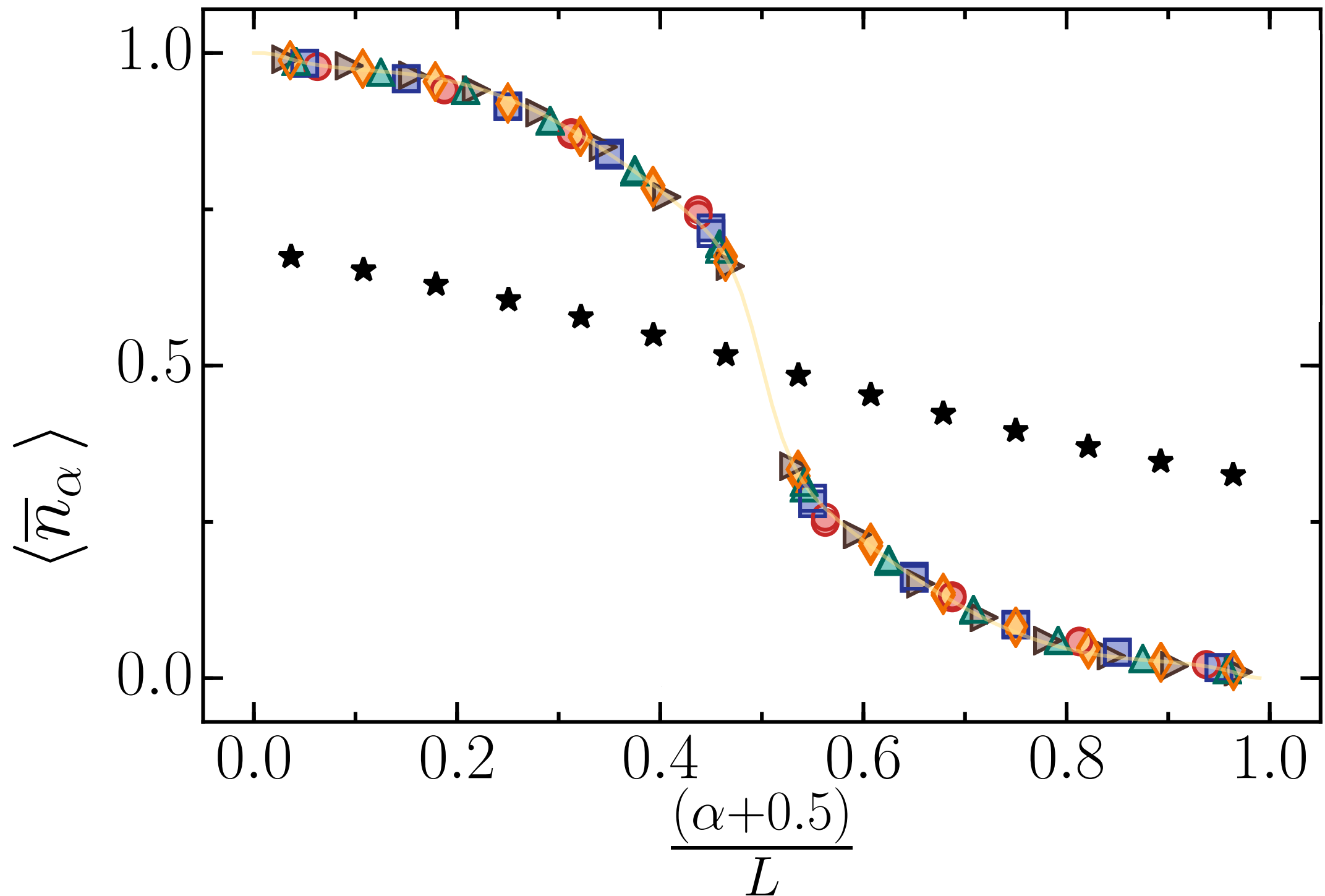
Initial correlations in CDW state destroyed by dephasing

$$\rho_{ij}(t) = \langle \psi(t) | c_i^\dagger c_j | \psi(t) \rangle$$

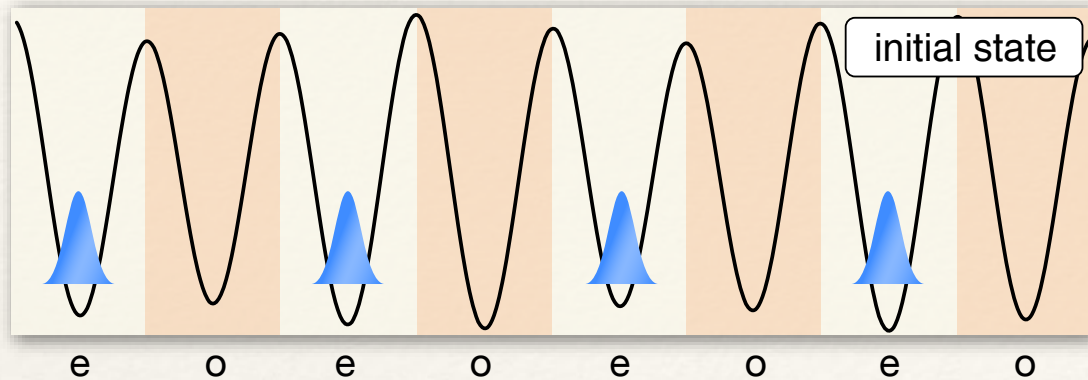
$$\rho(t)\phi_\alpha(t) = n_\alpha(t)\phi_\alpha(t)$$



Long time occupations nonthermal with smeared discontinuity

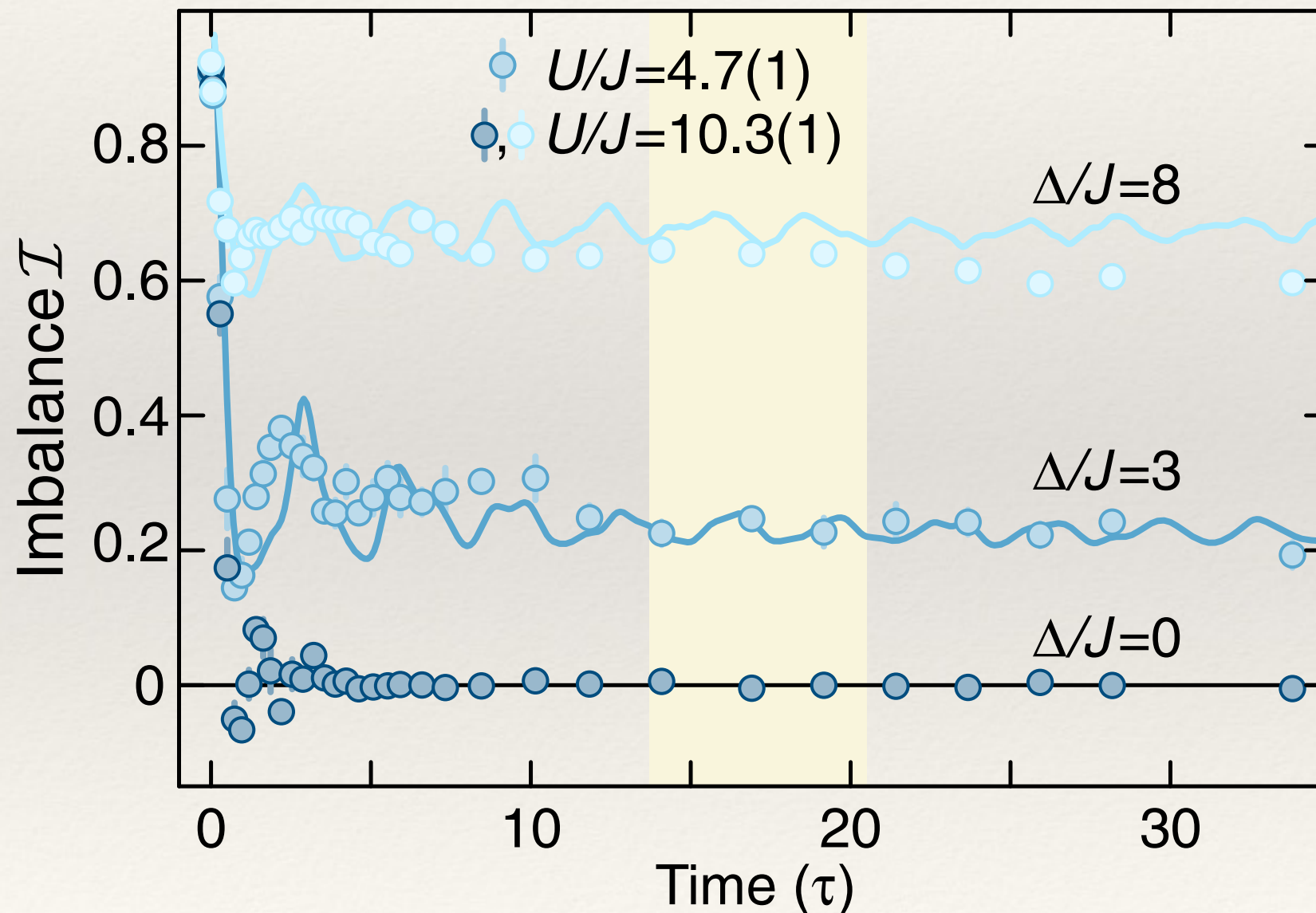


Many-body localization may be experimentally realized in cold atoms in a quasiperiodic lattice



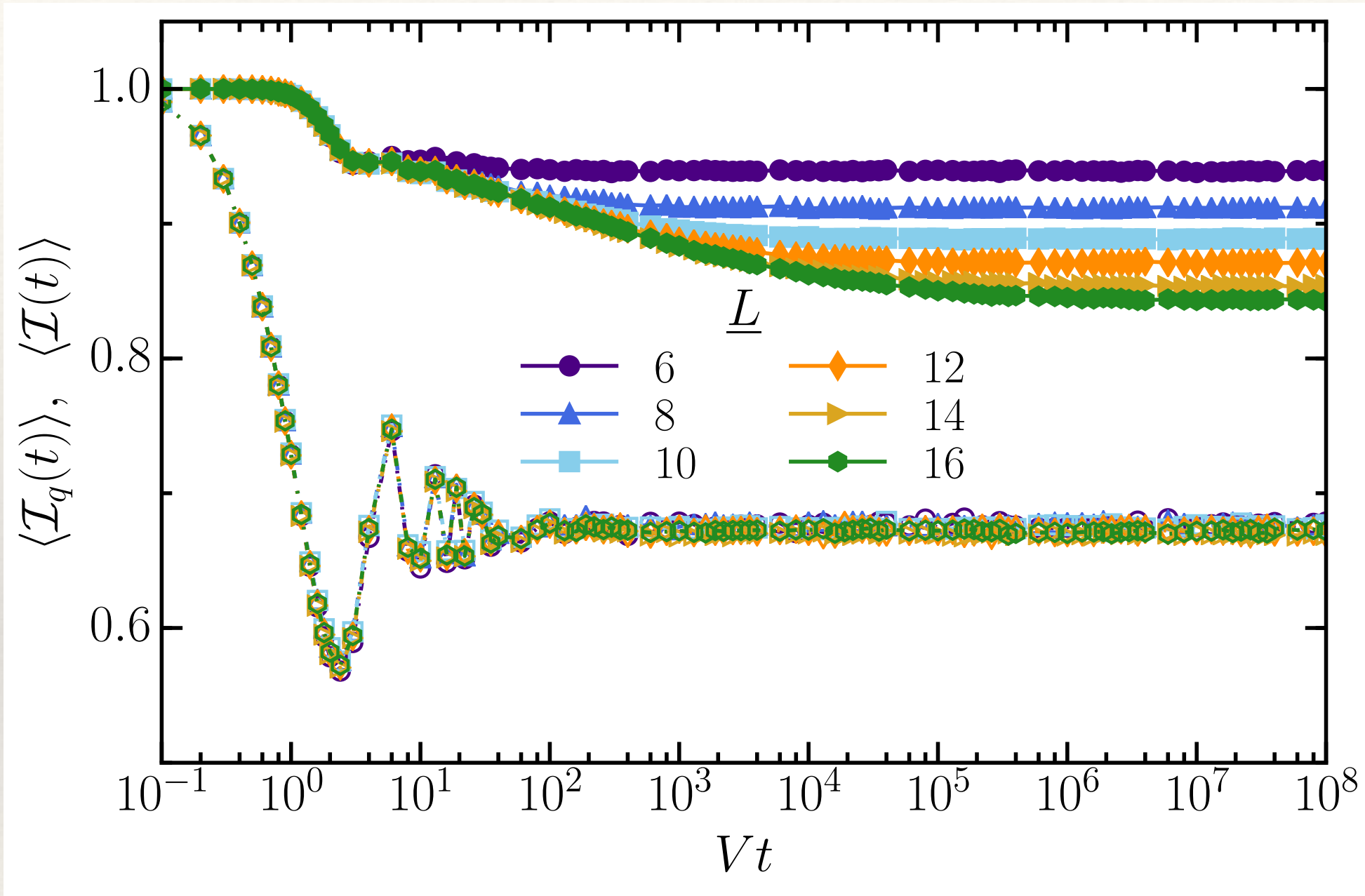
Aubry-André model

$$\hat{H} = -J \sum_{i,\sigma} (\hat{c}_{i,\sigma}^\dagger \hat{c}_{i+1,\sigma} + \text{h.c.}) + \Delta \sum_{i,\sigma} \cos(2\pi\beta i + \phi) \hat{c}_{i,\sigma}^\dagger \hat{c}_{i,\sigma} + U \sum_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$$



Occupation imbalance similar to density imbalance, but with slower relaxation towards steady state

$$\mathcal{I} = \frac{N_+ - N_-}{N}$$



Collaborators



Talía Lezama
Dresden



Giuseppe De Tomasi
Dresden



Soumya Bera
Dresden



Frank Pollmann
München



Henning Schomerus
Lancaster



Fabian Heidrich-Meisner
München

Summary

