## Multipartite Bell's Inequalities



Tomasz Paterek
Institute of Theoretical Physics and Astrophysics University of Gdańsk Poland

Collaborators:
Wieslaw Laskowski
Marek Zukowski
Caslav Brukner

## Bells Theorem

$\checkmark$ Quantum Probabilities
$\checkmark$ Underlying Theory
reproduces quantum probabilities as averages over inaccessible variables
$\checkmark$ Hidden Variable Theories hidden variables describe the properties of physical objects
$\checkmark$ Bell's Theorem it is impossible to explain all quantum mechanical expectations with local hidden-variable theories
A. Einstein, B. Podolsky, and N. Rosen, Phys. Rev. 47, 777 (1935).
$\square$ J. S. Bell, Physics 1, 195 (1964).

## Clauser-Horne-Shimony-Holt


$\checkmark$ realism
undone measurements have well-defined, yet unknown, results
$\checkmark$ locality
no "action at a distance"
$\checkmark$ For each experimental run:

$$
A_{1}\left[B_{1}+B_{2}\right]+A_{2}\left[B_{1}-B_{1}\right]= \pm 2
$$

$\checkmark$ Average over many runs:

$$
\begin{gathered}
-2 \leq\left\langle A_{1} B_{1}\right\rangle+\left\langle A_{1} B_{2}\right\rangle+\left\langle A_{2} B_{1}\right\rangle-\left\langle A_{2} B_{2}\right\rangle \leq 2 \\
\left|E_{11}+E_{12}+E_{21}-E_{22}\right| \leq 2
\end{gathered}
$$

[ J. F Clauser, M. A. Horne, A. Shimony, and R. A. Holt, Phys. Rev. Lett. 23, 880 (1969).

## Quantum Prediction

$\checkmark$ Take the state $\left.\left.\left\langle\psi^{-}\right\rangle=1 / \sqrt{2}[0\rangle 1\right\rangle-|1\rangle 0\right\rangle$
$\checkmark$ Quantum correlation function:

$$
E_{k l}{ }^{Q M}=-\cos \varphi_{k l}
$$

$\checkmark$ All measurements in one plane they are determined by the angle relative to some fixed axis

$$
\begin{array}{ll}
a_{1}=0 & b_{1}=\pi / 4 \\
a_{2}=\pi / 2 & b_{2}=-\pi / 4
\end{array}
$$

$\checkmark$ CHSH expression:

$$
\begin{aligned}
& \left|\boldsymbol{E}_{11}+\boldsymbol{E}_{12}+\boldsymbol{E}_{21}-\boldsymbol{E}_{22}\right| \leq 2 \\
& \cos (\pi / 4)+\cos (\pi / 4)+\cos (\pi / 4)-\cos (3 \pi / 4)=2 \sqrt{2}
\end{aligned}
$$

## Correlation Polytope

$$
\left(\begin{array}{l}
E_{11} \\
E_{12} \\
E_{21} \\
E_{22}
\end{array}\right)=\sum_{A_{1}, A_{2}, B_{1}, B_{2}= \pm 1} P\left(A_{1}, A_{2}, B_{1}, B_{2}\right)\binom{A_{1}}{A_{2}} \otimes\binom{B_{1}}{B_{2}}
$$

$\checkmark$ Bell's inequalities $=$ facets of the polytope
$\checkmark$ N\&S condition for local realistic model is a set of inequalities
$\checkmark$ For each experimental run:

$$
\sum_{s_{1}, s_{2}= \pm 1} S\left(s_{1}, s_{2}\right)\left(A_{1}+s_{1} A_{2}\right)\left(B_{1}+s_{2} B_{2}\right)= \pm 4
$$

$\checkmark$ Average over many runs:

$$
\sum_{s_{1}, s_{2}= \pm 1} S\left(s_{1}, s_{2}\right)\left[E_{11}+s_{2} E_{12}+s_{1} E_{21}+s_{1} s_{2} E_{22}\right] \leq 4
$$

$\checkmark$ There are as many inequalities as different sign functions
$\checkmark$ The complete set is equivalent to the single inequality:

$$
\sum_{s_{1}, s_{2}= \pm 1} \mid E_{11}+s_{2} E_{12}+s_{1} E_{21}+s_{1} s_{2} E_{22} \leq 4
$$

$\checkmark$ Works for arbitrary number of qubits and two measurement settings
H. Weinfurter and M. Zukowski, Phys. Rev. A 64, 10102(R) (2001).
$\square$ R. F. Werner and M. W. Wolf, Phys. Rev. A 64, 32112 (2001).

- M. Zukowski and C. Brukner, Phys. Rev. Lett. 88, 210401 (2002).


## Nonclassical States

$\checkmark$ General N-qubit state:

$$
\rho=\frac{1}{2^{N}} \sum_{x_{1}, \ldots, x_{N}=0}^{3} T_{x_{1} . \ldots x_{N}} \sigma_{x_{1}}^{1} \otimes \ldots \otimes \sigma_{x_{N}}^{N}
$$

$\checkmark$ Quantum correlation function:

$$
E_{k_{1} \ldots k_{N}}^{Q M}=\operatorname{Tr}\left[\rho\left(\vec{n}_{k_{1}} \cdot \vec{\sigma} \otimes \ldots \otimes \vec{n}_{k_{N}} \cdot \vec{\sigma}\right)\right]=\sum_{x_{1}, \ldots, x_{N} 1}^{3} T_{x_{1} 1 . x_{N}}\left(\vec{n}_{k_{1}}\right)_{x_{1}} \ldots\left(\vec{n}_{k_{N}}\right)_{x_{N}}
$$

$\checkmark$ Condition for the general inequality to hold:

$$
\max \sum_{x_{1}, \ldots, x_{N}=1}^{2} T_{x_{1} \ldots x_{N}}^{2} \leq 1
$$

maximization is taken over all local measurement directions
$\checkmark$ For more than 2 parties this condition is only the necessary one violation of it implies violation of the inequality
$\checkmark$ There are pure entangled states which do not violate the inequality

## More Settings

$\checkmark$ Three-particle $4 \times 4 \times 2$ case

$$
\begin{gathered}
A_{12, s^{\prime}} \equiv \sum_{s_{1}, s_{2}= \pm 1} S^{\prime}\left(s_{1}, s_{2}\right)\left(A_{1}+s_{1} A_{2}\right)\left(B_{1}+s_{2} B_{2}\right)= \pm 4 \\
A_{34, S^{\prime \prime}} \equiv \sum_{s_{1}, s_{2}= \pm 1} S^{\prime \prime}\left(s_{1}, s_{2}\right)\left(A_{3}+s_{1} A_{4}\right)\left(B_{3}+s_{2} B_{4}\right)= \pm 4 \\
\sum_{s_{1}, s_{2}= \pm 1} S\left(s_{1}, s_{2}\right)\left(A_{12, S^{\prime}}+s_{1} A_{34, S^{\prime \prime}}\right)\left(C_{1}+s_{2} C_{2}\right)= \pm 16
\end{gathered}
$$

$\checkmark$ There are $\left(2^{4}\right)^{3}$ inequalities
they are generated by one inequality with $\mathrm{S}, \mathrm{S}^{\prime}, \mathrm{S}^{\prime \prime}$ non-factorable
$\checkmark$ Incomplete set
$\checkmark$ Tight inequalities
[ W. Laskowski, TP, M. Zukowski, and C. Brukner, Phys. Rev. Lett. 93, 200401 (2004).
$\checkmark N \& S$ for the $4 \times 4 \times 2$ inequality to hold:

$$
\max \sum_{m=1}^{2} \sum_{k_{m}, l_{m}=1}^{2} T_{k_{m} l_{m} m}^{2} \leq 1
$$

$\checkmark$ Necessary for the $2 \times 2 \times 2$ case:

$$
\max \sum_{k, l, m=1}^{2} T_{k l m}^{2} \leq 1
$$

$\checkmark$ New condition is more demanding the settings $k_{1}, l_{1}$ do not have to be equal to $k_{2}, l_{2}$
$\checkmark$ The generalized GHZ states
violate new inequality for all alpha and arbitrary number of particles

$$
\psi\rangle=\cos \alpha|0 \ldots 0\rangle+\sin \alpha|1 \ldots 1\rangle
$$

