Multipartite Bell's Inequalities



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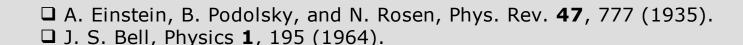
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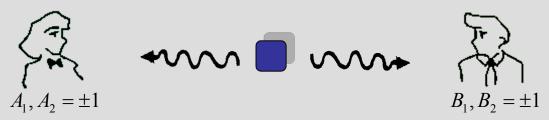
Bell's Theorem

- ✓ Quantum Probabilities
- ✓ Underlying Theory reproduces quantum probabilities as averages over inaccessible variables
- ✓ Hidden Variable Theories
 hidden variables describe the properties of physical objects
- ✓ Bell's Theorem
 it is impossible to explain all quantum mechanical expectations with local hidden-variable theories





Clauser-Horne-Shimony-Holt



- ✓ realism
 undone measurements have well-defined, yet unknown, results
- ✓ locality no "action at a distance"
- ✓ For each experimental run:

$$A_1[B_1 + B_2] + A_2[B_1 - B_1] = \pm 2$$

✓ Average over many runs:

$$-2 \le \langle A_1 B_1 \rangle + \langle A_1 B_2 \rangle + \langle A_2 B_1 \rangle - \langle A_2 B_2 \rangle \le 2$$
$$\left| E_{11} + E_{12} + E_{21} - E_{22} \right| \le 2$$

□ J. F Clauser, M. A. Horne, A. Shimony, and R. A. Holt, Phys. Rev. Lett. **23**, 880 (1969).

Quantum Prediction

- ✓ Take the state $|\psi^{-}\rangle = 1/\sqrt{2} [|0\rangle|1\rangle |1\rangle|0\rangle$
- ✓ Quantum correlation function:

$$E_{kl}^{QM} = -\cos\varphi_{kl}$$

✓ All measurements in one plane they are determined by the angle relative to some fixed axis

$$a_1 = 0$$
 $b_1 = \pi / 4$ $a_2 = \pi / 2$ $b_2 = -\pi / 4$

✓ CHSH expression:

$$\left| E_{11} + E_{12} + E_{21} - E_{22} \right| \le 2$$

$$\cos(\pi/4) + \cos(\pi/4) + \cos(\pi/4) - \cos(3\pi/4) = 2\sqrt{2}$$

Correlation Polytope

$$\begin{pmatrix} E_{11} \\ E_{12} \\ E_{21} \\ E_{22} \end{pmatrix} = \sum_{A_1, A_2, B_1, B_2 = \pm 1} P(A_1, A_2, B_1, B_2) \begin{pmatrix} A_1 \\ A_2 \end{pmatrix} \otimes \begin{pmatrix} B_1 \\ B_2 \end{pmatrix}$$

- ✓ Bell's inequalities = facets of the polytope
- √ N&S condition for local realistic model is a set of inequalities.

Two Qubits

✓ For each experimental run:

$$\sum_{s_1, s_2 = \pm 1} S(s_1, s_2) (A_1 + s_1 A_2) (B_1 + s_2 B_2) = \pm 4$$

✓ Average over many runs:

$$\left| \sum_{s_1, s_2 = \pm 1} S(s_1, s_2) [E_{11} + s_2 E_{12} + s_1 E_{21} + s_1 s_2 E_{22}] \right| \le 4$$

- ✓ There are as many inequalities as different sign functions
- ✓ The complete set is equivalent to the single inequality:

$$\sum_{s_1, s_2 = \pm 1} \left| E_{11} + s_2 E_{12} + s_1 E_{21} + s_1 s_2 E_{22} \right| \le 4$$

- ✓ Works for arbitrary number of qubits and two measurement settings
- ☐ H. Weinfurter and M. Zukowski, Phys. Rev. A **64**, 10102(R) (2001).
- ☐ R. F. Werner and M. W. Wolf, Phys. Rev. A **64**, 32112 (2001).
- ☐ M. Zukowski and C. Brukner, Phys. Rev. Lett. **88**, 210401 (2002).

Nonclassical States

✓ General N-qubit state:

$$\rho = \frac{1}{2^N} \sum_{x_1 = 0}^{3} T_{x_1 \dots x_N} \sigma_{x_1}^1 \otimes \dots \otimes \sigma_{x_N}^N$$

✓ Quantum correlation function:

$$E_{k_{1}...k_{N}}^{QM} = Tr[\rho(\vec{n}_{k_{1}} \cdot \vec{\sigma} \otimes ... \otimes \vec{n}_{k_{N}} \cdot \vec{\sigma})] = \sum_{x_{1},...,x_{N}=1}^{3} T_{x_{1}...x_{N}} (\vec{n}_{k_{1}})_{x_{1}}...(\vec{n}_{k_{N}})_{x_{N}}$$

✓ Condition for the general inequality to hold:

$$\max \sum_{x_1, \dots, x_N = 1}^{2} T_{x_1 \dots x_N}^2 \le 1$$

maximization is taken over all local measurement directions

- ✓ For more than 2 parties this condition is only the necessary one violation of it implies violation of the inequality
- ✓ There are pure entangled states which do not violate the inequality
- ☐ M. Zukowski, C. Brukner, W. Laskowski, and M. Wiesniak, Phys. Rev. Lett. 88, 210402 (2002)

More Settings

√ Three-particle 4x4x2 case

$$A_{12,S'} \equiv \sum_{s_1,s_2=\pm 1} S'(s_1,s_2)(A_1 + s_1 A_2)(B_1 + s_2 B_2) = \pm 4$$

$$A_{34,S''} \equiv \sum_{s_1,s_2=\pm 1} S''(s_1,s_2)(A_3 + s_1 A_4)(B_3 + s_2 B_4) = \pm 4$$

$$\sum_{s_1,s_2=\pm 1} S(s_1,s_2)(A_{12,S'} + s_1 A_{34,S''})(C_1 + s_2 C_2) = \pm 16$$

- ✓ There are (2⁴)³ inequalities they are generated by one inequality with S,S',S" non-factorable
- ✓ Incomplete set
- ✓ Tight inequalities

□ W. Laskowski, TP, M. Zukowski, and C. Brukner, Phys. Rev. Lett. 93, 200401 (2004).

More Nonclassical States

✓ N&S for the 4x4x2 inequality to hold:

$$\max \sum_{m=1}^{2} \sum_{k_m, l_m=1}^{2} T_{k_m l_m m}^2 \le 1$$

✓ Necessary for the 2x2x2 case:

$$\max \sum_{k,l,m=1}^{2} T_{klm}^2 \le 1$$

- ✓ New condition is more demanding the settings k_1 , l_1 do not have to be equal to k_2 , l_2
- ✓ The generalized GHZ states
 violate new inequality for all alpha and arbitrary number of particles

$$|\psi\rangle = \cos\alpha |0...0\rangle + \sin\alpha |1...1\rangle$$