

Suppression of fluorescence in a lossless cavity

A survey of the work by Alsing et al, PRA'92

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What we will do

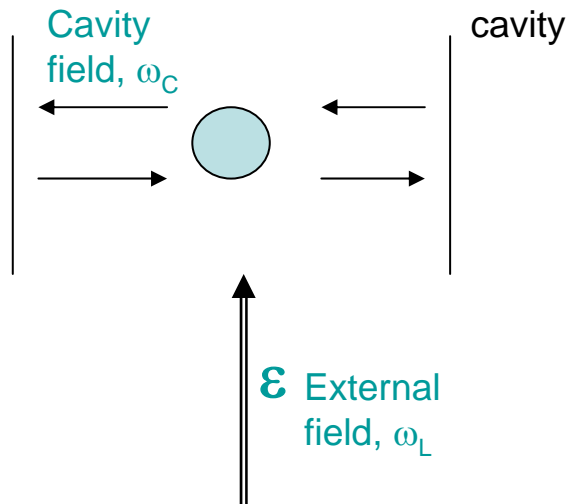
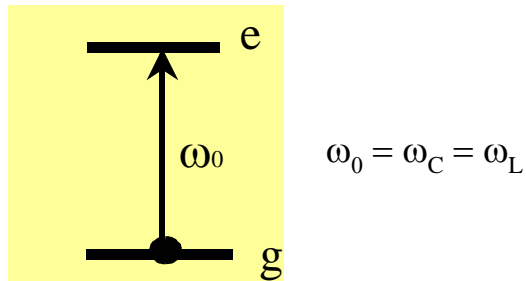
- We will review the paper by P.M. Alsing, D.A. Cardimona and H. J. Carmichael, Phys. Rev. A **45**, 1793 (1992) :

Investigation of the behaviour of a **two-level atom** interacting with both

- (i) a **single mode** of the e.m. field in a **lossless cavity** and
- (ii) an **external laser field**.

- We will discuss possible extensions of this work.

DESCRIPTION OF THE SYSTEM



- 1 two-level atom
- 2 single quantized cavity mode
- 3 Classical coherent external field

SYSTEM DYNAMICS

Interaction atom-cavity mode: Jaynes-Cummings Hamiltonian

$$\hat{H} = \underbrace{\frac{1}{2}\omega_0\hat{\sigma}_z}_{\text{atom}} + \underbrace{\omega_0\hat{a}^\dagger\hat{a}}_{\text{cavity}} + \underbrace{ig(\hat{a}^\dagger\hat{\sigma}_- - \hat{a}\hat{\sigma}_+)}_{\text{interaction}}$$

Master equation:

$$\begin{aligned} \frac{d\hat{\rho}}{dt} = & -i[\hat{H}, \hat{\rho}] + \underbrace{\varepsilon[\hat{\sigma}_-e^{i\omega_0t} - \hat{\sigma}_+e^{-i\omega_0t}, \hat{\rho}]}_{\text{interaction laser-atom}} \\ & + \underbrace{\frac{\gamma}{2}(2\hat{\sigma}_-\hat{\rho}\hat{\sigma}_+ - \hat{\sigma}_+\hat{\sigma}_-\hat{\rho} - \hat{\rho}\hat{\sigma}_+\hat{\sigma}_-)}_{\text{spontaneous emission}} \\ & + \underbrace{\kappa(2\hat{a}\hat{\rho}\hat{a}^\dagger - \hat{a}^\dagger\hat{a}\hat{\rho} - \hat{\rho}\hat{a}^\dagger\hat{a})}_{\text{cavity decay}} \end{aligned}$$

The atom is damped at rate γ by spontaneous emission

The cavity mode is damped at rate 2κ through the mirrors transmission

SEMICLASSICAL EQUATIONS

Mean intracavity field:

$$\alpha = \langle \hat{a} \rangle$$

$$\frac{d}{dt} \alpha = -\kappa \alpha + g(x) \nu$$

Atomic polarization:

$$\nu = \langle \sigma_- \rangle$$

$$\frac{d}{dt} \nu = -\frac{\gamma}{2} \nu - (g(x) \alpha + \varepsilon)$$

ATOMIC POLARIZATION

Steady-state solutions:
[below saturation]

In free space

$$v_f = -\frac{\varepsilon}{\gamma/2}$$

• Intensity of the fluorescence
proportional to $(2\varepsilon/\gamma)^2$

In the cavity

$$v_c = \frac{-\frac{\varepsilon}{\gamma/2}}{1 + \frac{2g^2}{\kappa\gamma}}$$

• Intensity proportional to

$$\left(\frac{\frac{2\varepsilon}{\gamma}}{1 + \frac{2g^2}{\kappa\gamma}} \right)^2$$

Modification
factor

$$\left(1 + \frac{2g^2}{\kappa\gamma} \right)^{-2}$$

SUPPRESSION OF FLUORESCENCE

$$|v_c|^2 = \left(\frac{\frac{2\varepsilon}{\gamma}}{1 + \frac{2g^2}{\kappa\gamma}} \right)^2$$

when

$$\frac{2g^2}{\kappa\gamma} \gg 1$$

then 

$$v_c = \frac{-\frac{\varepsilon}{\gamma/2}}{1 + \frac{2g^2}{\kappa\gamma}} \longrightarrow 0$$

THE CAVITY FIELD

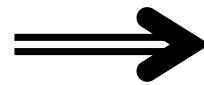
In the cavity, the steady-state solution of the mean intracavity field α is:

$$\alpha_c = \frac{g}{\kappa} \nu_c$$

and when

$$\frac{2g^2}{\kappa\gamma} \gg 1$$

then



$$\alpha_c \rightarrow -\varepsilon / g$$

And this result ($\alpha_c \rightarrow -\varepsilon / g$) means the situation where we have a suppression of the total field driving the atom:

$$\frac{d}{dt} \nu = -\frac{\gamma}{2} \nu - \underbrace{(g(x) \alpha + \varepsilon)}_{\substack{= \\ 0}}$$

SUMMARY

- One can achieve suppression of fluorescence in a cavity if

$$\frac{g^2}{\kappa\gamma} \gg 1 \quad (\text{large cooperativity})$$

- These results have been discussed in:

- 1 P.M. Alsing, D.A. Cardimona, H.J. Carmichael
Phys. Rev. A, **45**, 1793 (1992) [single atom, no motion, no spatial mode function]
- 2 S. Zippilli, G. Morigi, H. Ritsch, Phys. Rev. Lett. **93** (2004)
[generalization to spatial mode function, collective effects]

NEXT STEP

Atomic motion, spatial mode function

how does the cavity state changes?

$$\alpha_c = -\varepsilon / g(x)$$

$$v_c = \frac{-\frac{\varepsilon}{\gamma/2}}{1 + \frac{2g(x)^2}{\kappa\gamma}}$$

If $g(x) = g_0 \cos kx$ and $\cos kx = 0$

\implies This solution is NOT valid at the nodes of the cavity mode

- What happens if we move slowly (adiabatically) the atom from an antinode ($\cos kx = 1$) to a node ($\cos kx = 0$) ?