#### Calculating a Maximum Flux Transition Path

Robert D. Skeel, Purdue University, West Lafayette

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## Outline

#### Motivation

Objective

Calculation of Free Energy and Metric Tensor

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Discretization and Minimization

Weaknesses and Strengths of an MFTP

Footnotes and Closing Remarks

# **Motivation**



Catalytic domain of Src tyrosine kinase collaboration with Carol Post group

#### **Biological questions**

- What is the rate-limiting step in the transition from one conformation to another?
- What are possible intermediate states involved in the transition that can be used as targets for inhibitors of enhanced specificity?
- What is the free energy difference between the two conformational states?

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## Informal problem statement

Compute for minimal cost best possible representative paths of conformational change from conformation A to conformation B.

representative path: center of an isolated cluster of trajectories.

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#### Collective variables

Transition paths might not cluster adequately —in full configuration space. Assume, however, there is a smaller set of *collective variables*, functions of the configuration *x*,

$$\zeta_1 = \xi_1(x), \zeta_2 = \xi_2(x), \dots, \zeta_k = \xi_k(x),$$
 abbreviated as  $\zeta = \xi(x),$ 

such that in colvar space,

paths cluster into one or several distinct isolated bundles.

e.g., torsion angles  $\phi$  and  $\psi.$ 

Best possible: tube in colvar space of specified cross-section area for which the flow rate of distinct reactive trajectories is maximum —Best local maxima wanted.

Simplification: a *narrow* tube.

This formulation provides for comparisons among isolated bundles of paths.

Minimal cost: make simplifications to limit sampling to paths in colvar space. (Also, minimize programming effort by using existing features of simulators.)

-a maximum flux transition path (MFTP)

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#### Free energy function

Assume Newtonian dynamics with initial positions x and velocities drawn from the canonical ensemble with inverse temperature  $\beta$ . Let  $\rho_{\xi}(\zeta)$  be the probability density of  $\zeta = \xi(x)$ . An "effective energy" function  $F(\zeta)$  is obtained via

$$\mathrm{e}^{-\beta F(\zeta)} \stackrel{\mathrm{def}}{=} \rho_{\xi}(\zeta) = \langle \delta(\xi(x) - \zeta) \rangle,$$

sometimes called a free energy function.

#### Metric tensor

The appropriate metric to measure distance from  $\zeta$  to  $\zeta + d\zeta$  is

$$|\mathrm{d}\zeta|_{\zeta} = (\mathrm{d}\zeta^{\mathsf{T}}M(\zeta)\mathrm{d}\zeta)^{1/2}$$

with metric tensor

$$M(\zeta)^{-1} \stackrel{\text{def}}{=} m_{\text{tot}} \left\langle \xi_x(x) M_0^{-1} \xi_x(x)^{\mathsf{T}} \right\rangle_{\xi(x) = \zeta}$$

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where  $M_0$  is a diagonal matrix of masses.

#### The formula

The maximum flux transition path (MFTP)

$$\zeta = Z(s), \quad 0 \le s \le 1,$$

minimizes

$$\int_0^1 e^{\beta F(Z)} (\det M(Z))^{-1/2} \underbrace{|Z_s|_Z ds}_{\text{arc length}}$$

where Z = Z(s),  $Z_s = (d/ds)Z(s)$ , with  $Z(0) = \operatorname{argmin}_{\zeta \in A_{\xi}}F(\zeta)$  and  $Z(1) = \operatorname{argmin}_{\zeta \in B_{\xi}}F(\zeta)$ .

Zhao, Shen, and Skeel, *J. Chem. Theory Comput.*, 2010, building on transition path theory of E. Vanden-Eijnden and W. E.

Claim: An MFTP is the best representation of trajectories obtainable at minimal computing cost with modest programming effort.

at minimal computing cost with modest programming effort. Discussion deferred.

Objective: a more efficient algorithm for calculating an MFTP.

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Examples of MFTPs follow:

#### Three-hole potential



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The MFTP at different temperatures and the minimum energy path



MFEP vs. MFTP for alanine dipeptide in vacuo at T = 300 K. Contours of (zero-temperature) free energy in increments of 0.6 kcal/mol in  $\varphi$  and  $\psi$  torsion angles, red squares are MFTP, and black line is MFEP.

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#### Double basin Go model of CDK2 kinase

colvars  $\zeta = Z(s)$ 

Same final results for 3 different initial paths. Reorientation of  $\alpha$ -helix C is rate-limiting step.

free energy F(Z(s))



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#### The Dirac delta function

Convenient to work with a mollified delta function:

$$\delta_{arepsilon}(s) = (2\pi arepsilon^2)^{-1/2} \exp(-s^2/(2arepsilon^2))$$

e.g.,  $\varepsilon=1~{\rm degree}$ 

effectively stiff restraints

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### Energy function and metric tensor

Calculation of the free energy

$${\sf F}(\zeta) = -rac{1}{eta} \log \langle \delta(\xi(x)-\zeta) 
angle$$

requires extensive sampling.

However, the gradient can be expressed as an average conditioned on  $\xi(x) = \zeta$ :

$$abla F(\zeta) = -rac{1}{eta} \langle 
abla_{\zeta} \log \delta(\xi(x) - \zeta) 
angle_{\xi(x) = \zeta}.$$

Additionally, recall that

$$M(\zeta)^{-1} = m_{\text{tot}} \left\langle \xi_x(x) M_0^{-1} \xi_x(x)^{\mathsf{T}} \right\rangle_{\xi(x) = \zeta}$$

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## Averaging

It is practical to compute  $M(\zeta)$ ,  $\nabla_{\zeta}F(\zeta)$ , and 1st derivative of  $M(\zeta)$  w.r.t.  $\zeta$ 

as averages from a single (yet very long) simulation

at a point  $\zeta$  for which we have a value x such that  $\xi(x) \approx \zeta$  (to initiate the Markov chain).

an equilibration (burn-in) phase is followed by a production phase:

sampling for alanine dipeptide uses 50+500 ps of Langevin or Nosé-Hoover dynamics per gradient evaluation.

# $F(\zeta)$ and its Hessian

 Free energy differences can be constructed from free energy derivatives using piecewise quadratic interpolation.

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Calculating an approximate Hessian seems infeasible.

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## To simplify discussion

consider case of Cartesian coordinates scaled so that masses are 1: Find  $\zeta = Z(s)$ ,  $0 \le s \le 1$ , that minimizes

$$\int_0^1 \mathrm{e}^{\beta F(Z)} |Z_s| \mathrm{d}s$$

given Z(0) and Z(1).

(Difficulties due to colvars are technical only.)

# Gradient of the line integral $e^{\beta F(Z)} |Z_s| \left( I - \frac{Z_s Z_s^{\mathsf{T}}}{|Z_s|^2} \right) \left( \beta \nabla F(Z) - \frac{1}{|Z_s|^2} Z_{ss} \right)$

#### Challenges in discretizing

 arbitrariness in parameterization ζ = Z(s), manifested as a singularity in Euler-Lagrange equation

$$\left(I - \frac{Z_s Z_s^{\mathsf{T}}}{|Z_s|^2}\right) \left(-\frac{1}{|Z_s|^2} Z_{ss} + \beta \nabla F(Z)\right) = 0.$$

Possible normalization: s = relative arclength

$$|Z_s| = \text{ constant, i.e., } Z_s^{\mathsf{T}} Z_{ss} = 0.$$

... "benign" constraint.

#### • for large $\beta$ ,

the Euler-Lagrange equation is advection-dominated:

$$-\frac{1}{|Z_s|^2}Z_{ss}-\beta\frac{F(Z)_s}{|Z_s|^2}Z_s+\beta\nabla F(Z)=0.$$

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lack of energy function, very expensive gradient.

many minima We consider the computation of local minima (in path space).

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Start from an initial guess, which could be a straight line for a simple problem.

## Challenges in minimizing

- exponentially varying weights of terms in objective function
- Iack of convexity of free energy function.
- Iack of Hessian for free energy function.
- lack of free energy function, very expensive gradient.
- noisy gradient.
- cost of gradient evaluation increases with size of minimization step.

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#### Our current method

- piecewise linear discretization of path
- the semi-implicit simplified string method
   E, Ren, and Vanden-Eijnden. J. Chem. Phys., 2007.
   Vanden-Eijnden and Heymann. J. Chem. Phys., 2008.
  - (unprojected) gradient descent,

$$-\beta \nabla F(Z) + \frac{1}{|Z_s|^2} Z_{ss},$$

each step followed by reparameterization of the path.

equivalent to upwind differencing of Euler-Lagrange equations

Because most of the time is spent calculating  $\nabla F(\zeta)$  at replicas along the path, the computation is highly parallelizable:

MPI for Python: mpi4py

## String method drawbacks

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## 80 K temperature



## 79 K temperature



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discontinuous behavior for paths with discontinuous tangents.

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- lack of a readily computable line integral (a check for descent, a way to compare isolated paths)
- extra baggage of a dynamical embedding

## Direct approach to minimization

Discretizing the integral instead of the Euler-Lagrange equation provides search directions guaranteed to decrease the integral.

Let  $\varphi(z)$  be the discrete line integral where  $z = \begin{bmatrix} z_1^T & z_2^T & \dots & z_{J-1}^T \end{bmatrix}^T$  are the replica colvars defining the path.

#### A trust region approach

 $\varphi(z)$  not convex & partial Hessian available  $\Rightarrow$ consider a trust region approach. Such a region contains the current iterate  $z^{(k)}$ and represents the extent of our trust in a quadratic local model  $\tilde{\varphi}(z^{(k)} + w)$ for the objective function  $\varphi(z^{(k)} + w)$ .

Radius of region adjusted depending on  $\tilde{\varphi}(z^{(k+1)}) - \varphi(z^{(k+1)})$ .

Many challenges were seemingly overcome.

However, adjusting the trust region based on a local quadratic model does not work. What was I thinking!

## March 20, 2012

Local quadratic model is flawed. There is no tractable local model, it seems.

Conclusion: embrace a dynamical embedding for minimization which I have hitherto shunned.

#### Dynamics for minimization

Define a dynamical path  $\zeta = Z(s; t)$  by

$$\frac{\mathrm{d}}{\mathrm{d}t}Z = -S(Z)\nabla\varphi(Z)$$

where S(Z) is a diagonal scaling matrix.

- Scaling overcomes exponential range of weighting.
- Decrease of  $\varphi(Z)$  can be assured.
- Partial Hessian can be exploited by semi-implicit time stepping.
- Constraints readily accommodated, in principle.

Discretize first in space or in time?

Discretizing in space first seems more robust.

## March 25, 2012

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Observed apparent failure of line integral discretization, combined with normalization constraints. So far unable to revive. Have not given up.

Conceptually, the string method discretizes first in time.

## Outline

Motivation Objective Calculation of Free Energy and Metric Tensor Discretization and Minimization Weaknesses and Strengths of an MFTP Footnotes and Closing Remarks

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## Underlying assumption #1

Assume trajectories are *confined* to the tube from  $A_{\xi}$  to  $B_{\xi}$ .

This transforms a PDE in high dimension to one that is essentially one dimensional.

## Underlying assumption #2

Assume that paths  $\zeta = Z(s)$ , but not trajectories  $\zeta = \xi(X(t))$ , are well approximated by those of the Brownian dynamics.

#### For comparison

Minimum resistance path (MRP) minimizes

$$\int_0^1 \exp(\beta F(Z)) |Z_s|^{-1} |Z_s|_Z^2 \,\mathrm{d}s.$$

Path depends on how colvar space is parameterized (see The geometric property).

Berkowitz, Morgan, McCammon, and Northrup, J. Chem. Phys., 1983.Huo and Straub, "MaxFlux ...," J. Chem. Phys., 1997.

#### The geometric property

Use of metric  $|d\zeta|_{\zeta}$  to measure arc length and cross-section area. ensures that the minimizing path is invariant under a change of coordinates:

If we minimize the integral using instead variables

 $\zeta' = \chi(\xi(x)),$ 

the resulting path  $\zeta' = Z'(s)$  satisfies  $Z'(s) = \chi(Z(s))$ . e.g., squaring colvars yields same path. Minimum free energy path (MFEP) minimizes

$$\int_0^1 |M(Z)^{-1}\nabla F(Z)|_Z |Z_s|_Z \,\mathrm{d}s.$$

It is the limiting case  $\beta \to \infty$  of an MFTP but with  $\beta$  held fixed in  $F(\zeta; \beta)$  and  $M(\zeta; \beta)$ .

Maragliano, Fischer, Vanden-Eijnden, and Ciccotti, J. Chem. Phys., 2006.Vanden-Eijnden and Heymann. J. Chem. Phys., 2008. The MFEP neglects finite temperature effects —in the explicit degrees of freedom. alanine dipeptide (OPLS-AA with GBSA)



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Jiménez and Crehuet, Theor. Chem. Account, 2007.

## Outline

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## Initiating the Markov chains

Evaluation of  $\nabla F(Z_j)$  at time step n + 1requires  $x_j$  such that  $\xi(x_j) \approx Z_j$ . One can use for  $x_j$  a value generated during the production phase of sampling for  $Z_j$  at time step n, e.g., the last value. The number of equilibration steps depends on the change in  $Z_j$ . Greater increments result in longer equilibration times. Therefore, taking larger steps may or may not be helpful.

## A reference path in Cartesian space

In addition to  $\zeta = \xi(x)$ , we should define a *reference path* x = X(s) in Cartesian coordinates.

An attractive choice:

Ask that X(s) be an MFTP in Cartesian configuration space subject to the constraint  $\xi(X(s)) = Z(s)$ .

## Interpretation of $M(\zeta)$

The point  $\zeta$  in colvar space

represents a manifold in Cartesian configuration space. The distance

$$|\mathrm{d}\zeta|_{\zeta} = \left(\mathrm{d}\zeta^{\mathsf{T}}\left(\left\langle\xi_{x}(x)M_{0}^{-1}\xi_{x}(x)^{\mathsf{T}}\right\rangle_{\zeta}\right)^{-1}\mathrm{d}\zeta\right)^{1/2}$$

somehow corresponds to the RMSD between Cartesian space manifolds  $\xi(x) = \zeta$  and  $\xi(x) = \zeta + d\zeta$ 

$$\left(\mathrm{d}\zeta^{\mathsf{T}}\left\langle \left(\xi_{x}(x)M_{0}^{-1}\xi_{x}(x)^{\mathsf{T}}\right)^{-1}\right\rangle _{\zeta}\mathrm{d}\zeta\right)^{1/2}$$

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# **Closing Remarks**

To minimize an integral of an exponential, embed the problem in dynamics, as in the string method.

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 Yet significant improvements are desirable and almost certainly attainable.

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