

Statistical analysis and stochastic modelling of cell migration and bumblebee foraging

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Advanced Study Group, Second Focus Week Meeting
6 November 2015



Outline

two parts:

- 1 **cell migration**
- 2 **bumblebee foraging**

in both cases:

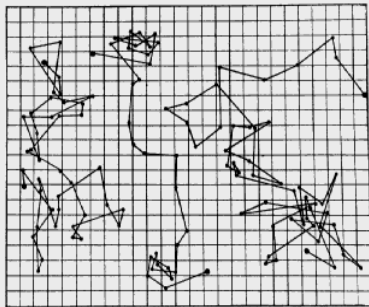
- **motivation and experiment**
- **experimental results and statistical analysis**
- **theoretical stochastic modeling and summary**

Part 1:

Cell Migration

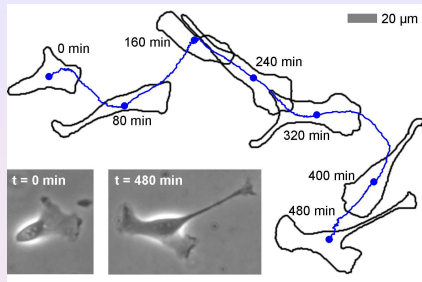
Brownian motion of migrating cells?

Brownian motion



Perrin (1913)

three colloidal particles,
positions joined by straight
lines



Dieterich et al. (2008)

single biological cell crawling on
a substrate

Brownian motion?

conflicting results:

yes: Dunn, Brown (1987)

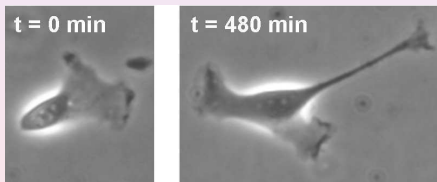
no: Hartmann et al. (1994)

Our cell types and how they migrate

MDCK-F (Madin-Darby canine kidney) cells

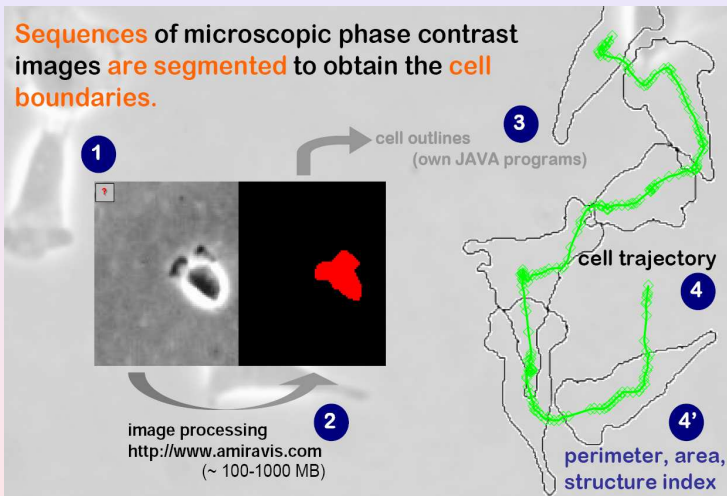
two types: wildtype (NHE^+) and NHE-deficient (NHE^-)

movies: NHE^+ : t=210min, dt=3min | NHE^- : t=171min, dt=1min



Measuring cell migration

Sequences of microscopic phase contrast images **are segmented** to obtain the **cell boundaries**.



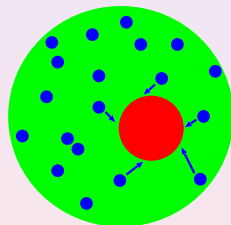
Theoretical modeling of Brownian motion

‘Newton’s law of stochastic physics’:

$$\dot{\mathbf{v}} = -\kappa\mathbf{v} + \sqrt{\zeta} \boldsymbol{\xi}(t) \quad \text{Langevin equation (1908)}$$

for a **tracer particle of velocity \mathbf{v}** immersed in a fluid

force decomposed into **viscous damping** and **random kicks of surrounding particles**

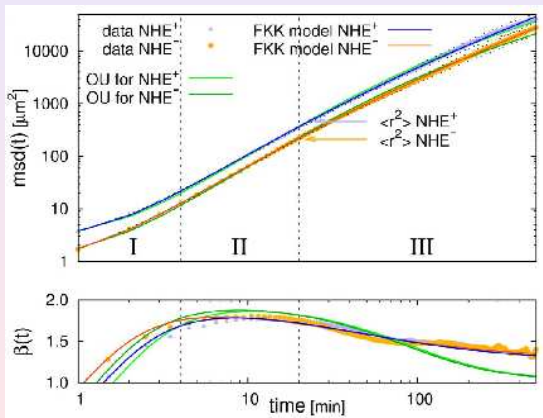


Application to cell migration?

but: cell migration is **active** motion, **not passively** driven!

Mean square displacement

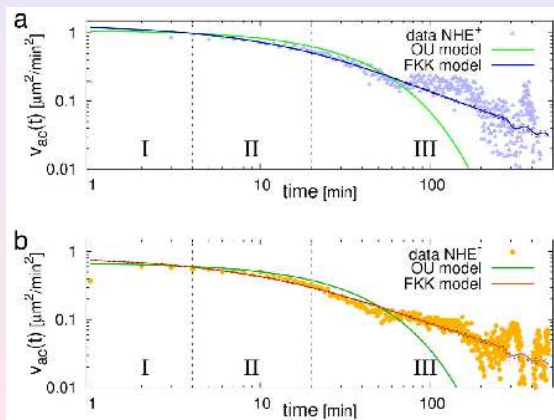
- $msd(t) := \langle [\mathbf{x}(t) - \mathbf{x}(0)]^2 \rangle \sim t^\beta$ with $\beta \rightarrow 2$ ($t \rightarrow 0$) and $\beta \rightarrow 1$ ($t \rightarrow \infty$) for Brownian motion; $\beta(t) = d \ln msd(t) / d \ln t$



anomalous diffusion if $\beta \neq 1$ ($t \rightarrow \infty$); here: **superdiffusion**

Velocity autocorrelation function

- $v_{ac}(t) := \langle \mathbf{v}(t) \cdot \mathbf{v}(0) \rangle \sim \exp(-\kappa t)$ for Brownian motion
- fits with same parameter values as $msd(t)$



crossover from **stretched exponential to power law**

Position distribution function

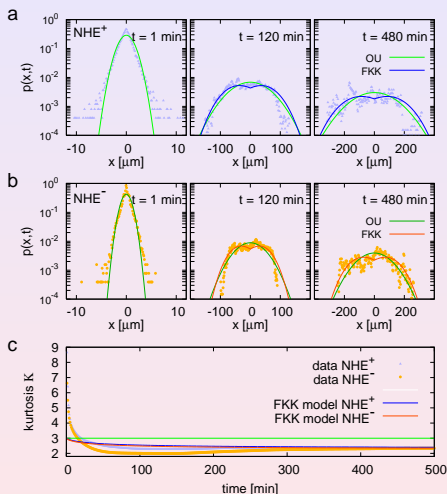
- $P(x, t) \rightarrow$ Gaussian ($t \rightarrow \infty$) and kurtosis

$$\kappa(t) := \frac{\langle x^4(t) \rangle}{\langle x^2(t) \rangle^2} \rightarrow 3 \quad (t \rightarrow \infty)$$

for Brownian motion (green lines, in 1d)

- other solid lines: fits from our model; parameter values as before

note: model needs to be amended to explain short-time distributions



crossover from peaked to broad **non-Gaussian distributions**

The model

- **Fractional Klein-Kramers equation** (Barkai, Silbey, 2000):

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x} [vP] + \frac{\partial^{1-\alpha}}{\partial t^{1-\alpha}} \kappa \left[\frac{\partial}{\partial v} v + v_{th}^2 \frac{\partial^2}{\partial v^2} \right] P$$

with probability distribution $P = P(x, v, t)$, damping term κ , thermal velocity $v_{th}^2 = kT/m$ and **Riemann-Liouville fractional** (generalized ordinary) **derivative of order $1 - \alpha$**

for $\alpha = 1$ Langevin's theory of Brownian motion recovered

- **analytical solutions** for $msd(t)$ and $P(x, t)$ can be obtained in terms of special functions (Barkai, Silbey, 2000; Schneider, Wyss, 1989)

- **4 fit parameters** v_{th}, α, κ (plus another one for short-time dynamics)

Possible physical interpretation

Physical meaning of the fractional derivative?

the **generalized Langevin equation**

$$\dot{v} + \int_0^t dt' \kappa(t-t')v(t') = \sqrt{\zeta} \xi(t)$$

e.g., Mori, Kubo (1965/66)

with **time-dependent friction coefficient** $\kappa(t) \sim t^{-\alpha}$ generates *the same* $msd(t)$ and $v_{ac}(t)$ as the fractional Klein-Kramers equation

cell anomalies might originate from **glassy behavior** of the cytoskeleton gel, where power law exponents are conjectured to be universal (Fabry et al., 2003; Kroy et al., 2008)

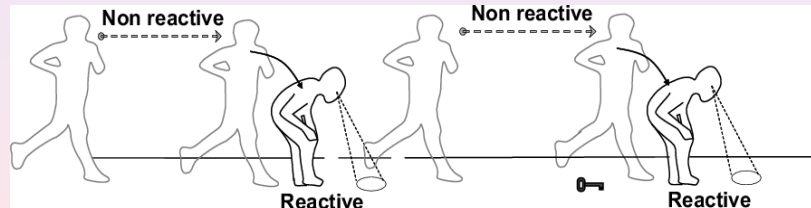
nb: anomalous dynamics observed for *many different cell types*

Possible biological interpretation

Biological meaning of the anomalous cell migration?

experimental data and theoretical modeling suggest *slower diffusion for small times* while *long-time motion is faster*

compare with **intermittent optimal search strategies** of foraging animals (Bénichou et al., 2006)



note: controversy about **modeling the migration of foraging animals** (albatros, **bumblebees**, fruitflies,...)

Summary: Anomalous cells

- different **cell dynamics** on different **time scales**
(cp. with **Lévy hypothesis**, which suggests scale-freeness)
- for long times cells crawl **superdiffusively** with **power law decay of velocity correlations and non-Gaussian position pdfs**
- **stochastic modeling** of experimental data by a **generalized Klein-Kramers equation**

Part 2:

Bumblebee Foraging

Motivation

bumblebee foraging – two very practical problems:

1. find food (nectar, pollen) in complex landscapes



2. try to avoid predators

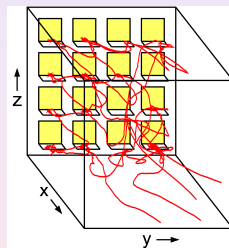
What type of motion?

Study bumblebee foraging in a *laboratory experiment*.

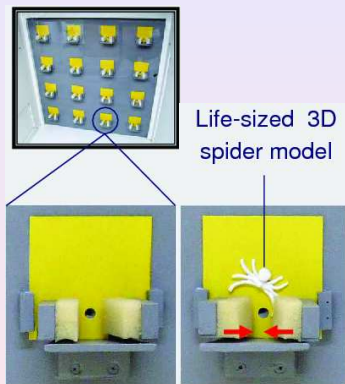
The bumblebee experiment

Ings, Chittka, *Current Biology* **18**, 1520 (2008):
bumblebee foraging in a cube of $\simeq 75\text{cm}$ side length

- artificial yellow flowers: 4x4 grid on one wall
- two cameras track the position (50fps) of a single bumblebee (*Bombus terrestris*)
- **advantages:** systematic **variation of the environment**;
easier than tracking bumblebees on large scales
- **disadvantage:** no 'free flight' of bumblebees



Variation of the environmental conditions



safe and **dangerous**
flowers

movie

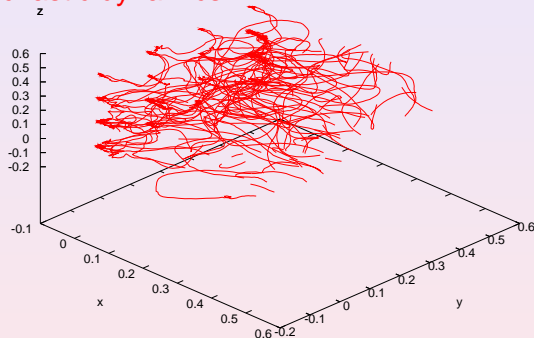
three experimental stages:

- 1 spider-free foraging
- 2 foraging under predation risk
- 3 memory test 1 day later

#bumblebees=30 , #data per bumblebee for each stage \approx 7000

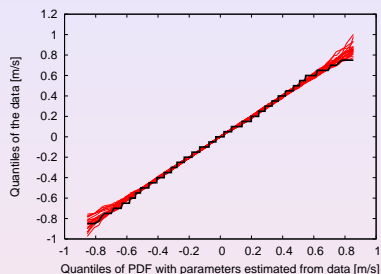
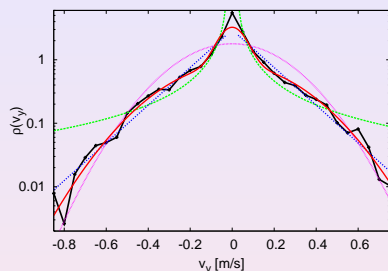
Bumblebee experiment: two main questions

- 1 What **type of motion** do the bumblebees perform in terms of **stochastic dynamics**?



- 2 Are there **changes of the dynamics** under **variation of the environmental conditions**?

Velocity distributions: analysis



left: experimental **pdf of v_y -velocities** of a single bumblebee in the spider-free stage (black crosses) with max. likelihood fits of **mixture of 2 Gaussians**; **exponential**; **power law**; **single Gaussian**

right: **quantile-quantile plot** of a Gaussian mixture against the experimental data (black) plus **surrogate data**

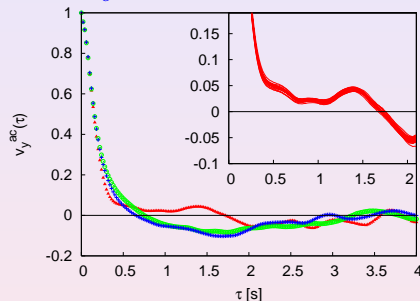
Velocity distributions: interpretation

- **best fit** to the data by a **mixture of two Gaussians** with different variances (quantified by information criteria with resp. weights)
- **biological explanation:** models **spatially different flight modes** near the flower vs. far away, cf. intermittent dynamics

big surprise: no difference in pdf's between different stages under variation of environmental conditions!

Velocity autocorrelation function || to the wall

$$V_y^{AC}(\tau) = \frac{\langle (v_y(t) - \mu)(v_y(t+\tau) - \mu) \rangle}{\sigma^2} \text{ with average over all bees:}$$

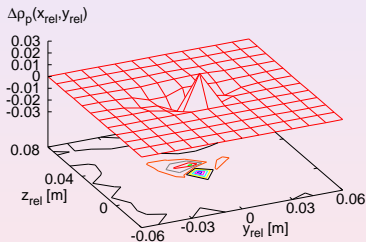


- plot: spider-free stage, predation thread, memory test
- **correlations change** from positive (spider-free) to negative (spiders)

⇒ all **changes** are in the **velocity correlations**, *not* in pdf's!

Predator avoidance and a simple model

predator avoidance as
difference in position pdfs
spider / no spider from data:



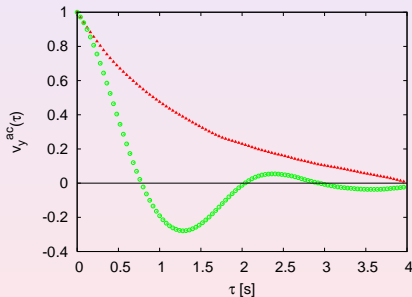
positive spike: *hovering*;
negative region: *avoidance*

modeled by Langevin equation

$$\frac{dv_y}{dt}(t) = -\eta v_y(t) - \frac{\partial U}{\partial y}(y(t)) + \xi(t)$$

η : friction coefficient,

ξ : Gaussian white noise



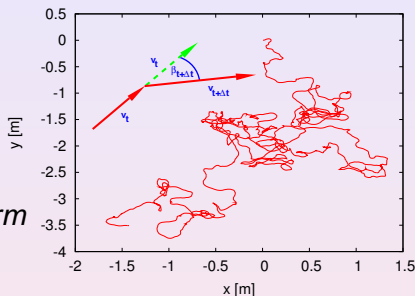
simulated velocity correlations with
repulsive interaction potential U
bumblebee - spider **off** / **on**

Modeling free bumblebee flights

reorientation model:

describe 2d movement in comoving frame by

- speed $v(t) = \text{const.}$
- turning angle $\beta(t) = \xi(t)$ as random variable from *non-uniform pdf* modeling **persistence**



generalized model for bumblebee flights far away from flowers constructed from experimental data:

- $\beta(t) = \xi_v(t)$: power law **correlated Gaussian noise**
- $\frac{dv}{dt} = g(v(t)) + \psi(t)$: **generalized Langevin equation** with anti-correlated Gaussian noise

Summary: Clever bumblebees

- mixture of **two Gaussian velocity distributions** reflects **spatial adjustment** of bumblebee dynamics to flower carpet
- all changes to predation threat are contained in the **velocity autocorrelation functions**, which exhibit highly **non-trivial temporal behaviour**
(nb: **Lévy hypothesis** suggests that all relevant foraging information is contained in pdf's)
- **change of correlation decay** in the presence of spiders due to **experimentally extracted repulsive force** as reproduced by generalized Langevin dynamics

Collaborators and literature

work performed with:

1. cells:

P.Dieterich, R.K., R.Preuss, A.Schwab, PNAS **105**, 459 (2008)

2. bees:

F.Lenz, T.Ings, A.V.Chechkin, L.Chittka, R.K.,
Phys. Rev. Lett. **108**, 098103 (2012)

F.Lenz, A.V.Chechkin, R.K., PLoS ONE 8, e59036 (2013)

